

TRANSITIONS IN SOME STAGNATION FLOWS OF VISCOELASTIC FLUIDS AT LOW REYNOLDS NUMBERS

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Flow Instabilities and Turbulence in Viscoelastic Fluids
Leiden, Holland July 19-23, 2010

OUTLINE

- Objective
- Experimental and numerical results
 - Cross slot
 - Flow focusing
- Some analytical thoughts: Stagnation + “vortex” flow
- Closure

OBJECTIVE

- **Elastic instabilities** ($Re=0$): enhanced mixing or upper limit in devices
- Transition from steady symmetric to steady asymmetric flow is our main interest
- When it occurs and what are the effects of solvent, inertia and extensional viscosity. Brief review in some simple flows
- Some findings about the asymmetric flow: decoupling into simpler flows
- Results: mostly numerical (some experiments) and analytical (work in

REVIEW

Viscoelastic instabilities in shear flows

Shaqfeh, Ann. Rev. Fluid Mech 28 (1996) 129

Taylor-Couette flow Larson et al., JFM 218 (1990) 573

Cone-plate flow McKinley et al., JNNFM 40 (1991) 201

Lid driven cavity flows Pakdel & McKinley, PRL 77 (1996) 2459

Underlying mechanism McKinley et al, JNNFM 67 (1996) 19

Pakdel & McKinley, PRL 77 (1996) 2459

$$\left(\frac{\lambda U}{\mathcal{R}} \frac{\tau_{11}}{\tau_{12}} \right) \geq M_{crit}^2 \quad \text{curved streamlines}$$

Instability growth to elastic turbulence

Groisman & Steinberg, Nature 405 (2000) 53 Larson, Nature 405 (2000) 27

Microfluidics & viscoelasticity

Squires & Quake, Rev. Mod. Phys. 77 (2005) 977

Transitions in some stagnation viscoelastic flows at Re=0
Flow Instabilities and turbulence in viscoelastic fluids

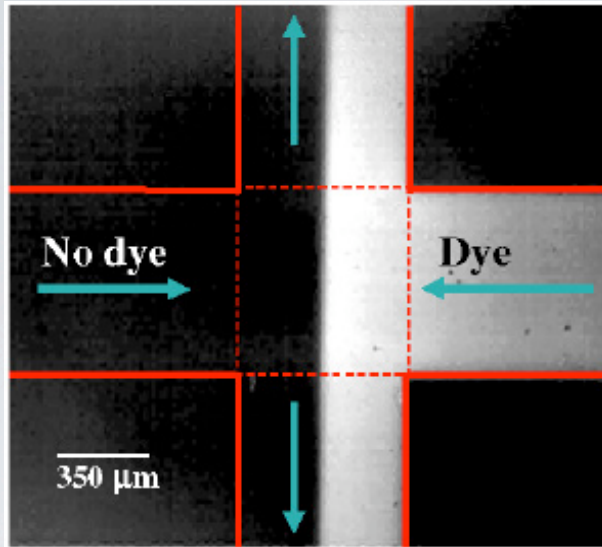
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Leiden, Holland, 19-23th July 2010

NUMERICAL METHODS: SOLUTION OF THE GOVERNING EQUATIONS

- Finite-volume method (in-house code)
- Collocated block-structured mesh
- Non-orthogonal coordinates (Cartesian velocity and stress tensor)
- Diffusion: central differences (2nd order in uniform mesh)
- SIMPLEC algorithm
- Rhie-and-Chow to couple velocity and pressure
- Special scheme to couple velocity and extra stress
[Oliveira et al. JNNFM, 79 \(1998\) 1-43.](#)
- Advection: CUBISTA high-resolution scheme (based on QUICK, 3rd order)
[Alves et al. IJNMF, 41 \(2003\) 47-75.](#)
- Standard formulation and log-conformation formulation (allows higher De)
[Fattal & Kupferman JNNFM, 123 \(2004\) 281-285.](#)
More details for FVM: [Afonso et al. JNNFM 157 \(2009\) 55-65](#)

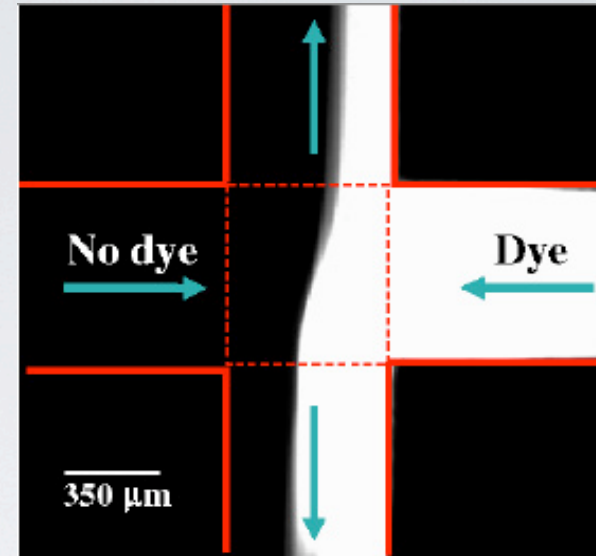
CROSS SLOT

2D CROSS SLOT

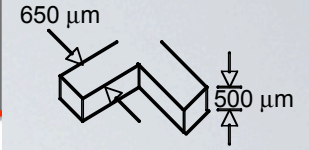


Newtonian ($Re < 10^{-2}$)

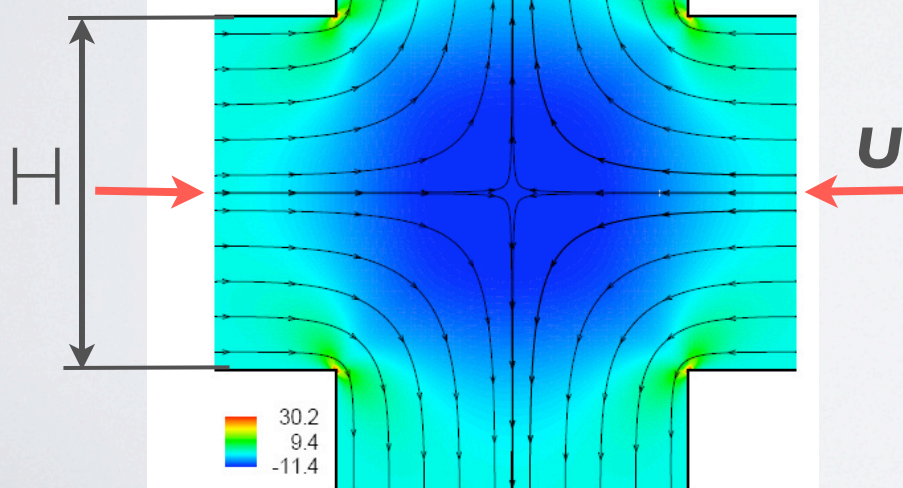
Arratia et al., PRL 96 (2006) 144502



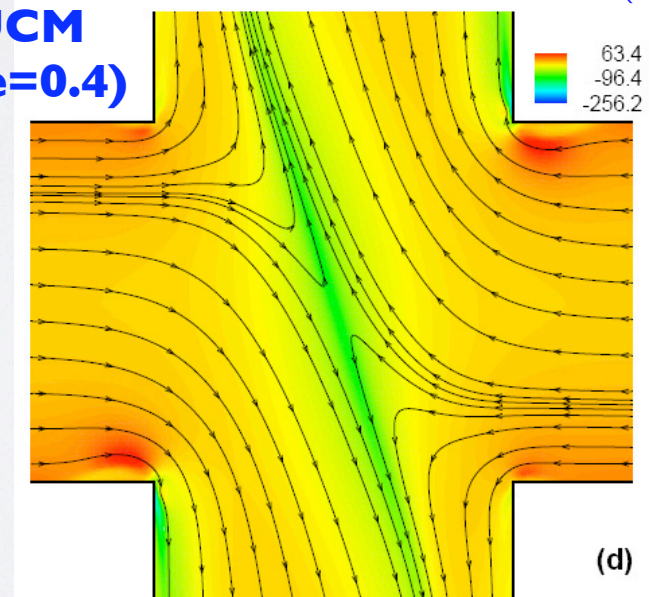
PAA Boger fluid ($Re < 10^{-2}$, $De=4.5$)



**Newtonian
($Re = 0$)**



**UCM
($De=0.4$)**

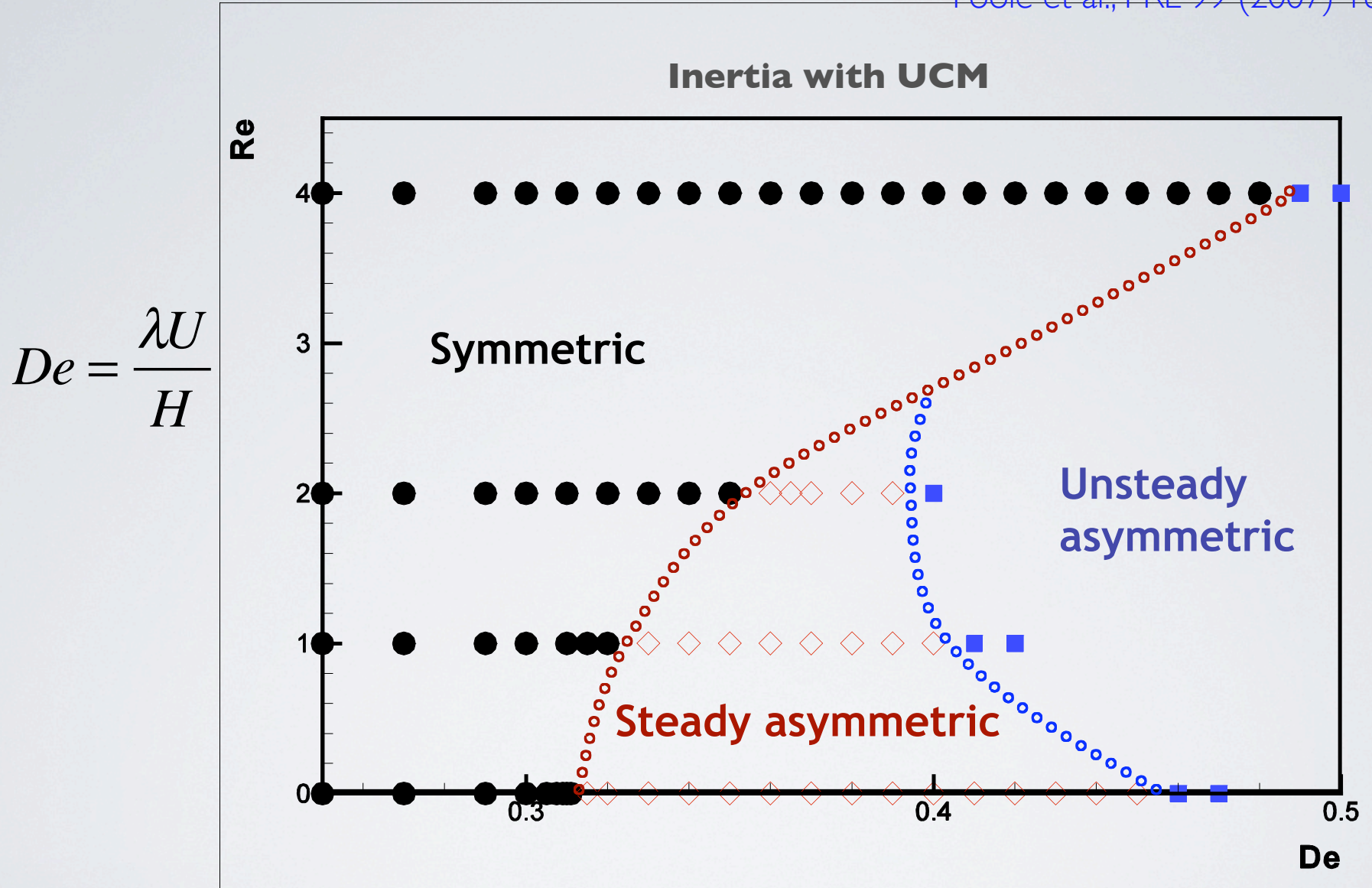


Transitions in some stagnation viscoelastic flows at $Re=0$
Flow Instabilities and turbulence in viscoelastic fluids

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2D CROSS SLOT WITH UCM: EFFECT OF INERTIA

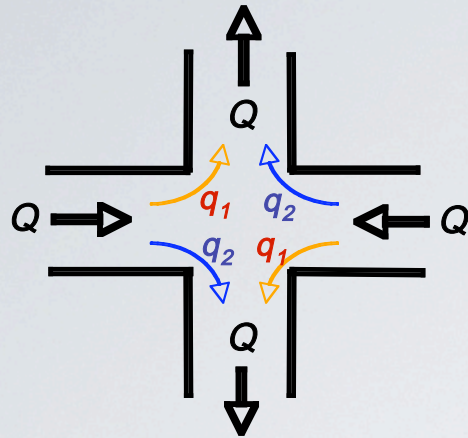
Poole et al., PRL 99 (2007) 164503



Inertia decreases degree of asymmetry and stabilizes the flow

2D CROSS SLOT: OLDROYD-B — EFFECT OF SOLVENT — CREEPING FLOW

Poole et al., SoR 2007



Oldroyd-B

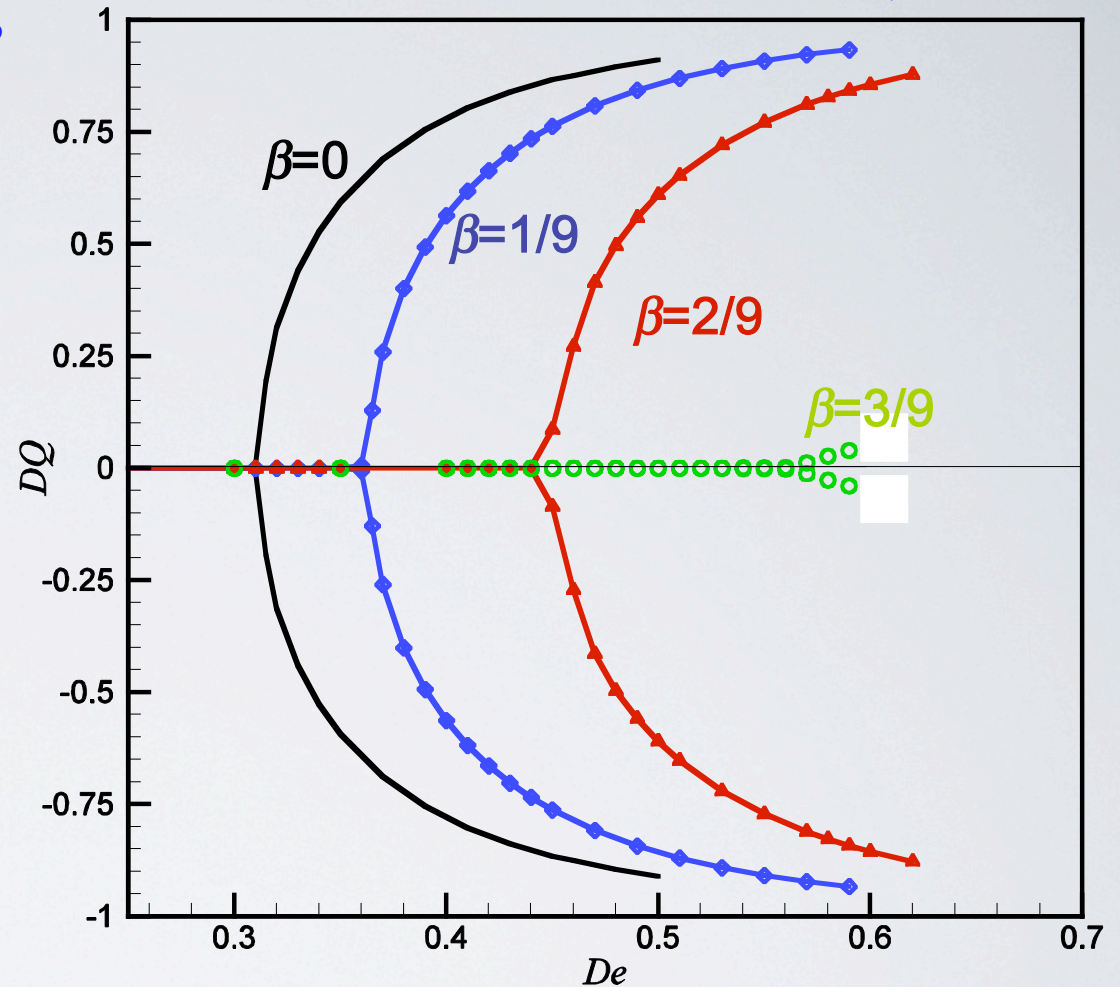
$Re=0$

$$DQ = \frac{q_2 - q_1}{q_2 + q_1} = \frac{q_2 - q_1}{Q}$$

$DQ = 0 \rightarrow$ symmetric

$DQ = \pm 1 \rightarrow$ completely asymmetric

$$\beta = \frac{\eta_s}{\eta_s + \eta_p}$$



Increasing the solvent viscosity

Increases De_{CR}

For $\beta > 3/9$ flow becomes asymmetric unsteady (as in flow focusing)

2D CROSS SLOT: OLDROYD-B — SOLVENT AND INERTIA

Poole et al., SoR 2007

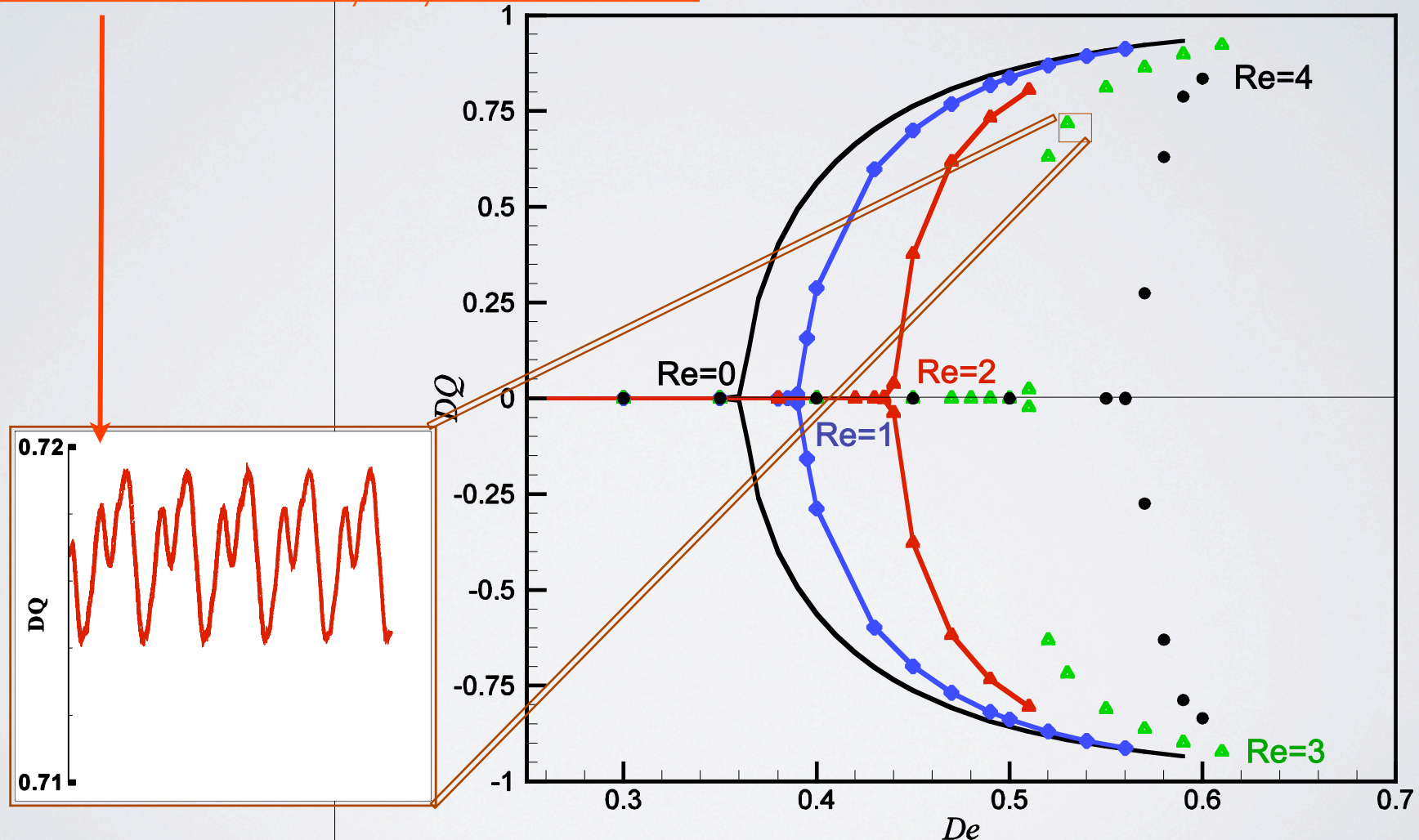
Increasing Re

$$\beta = 1/9$$

Increases De_{CR}

Decreases degree of asymmetry

For $Re > 2$ unsteady asymmetric flow

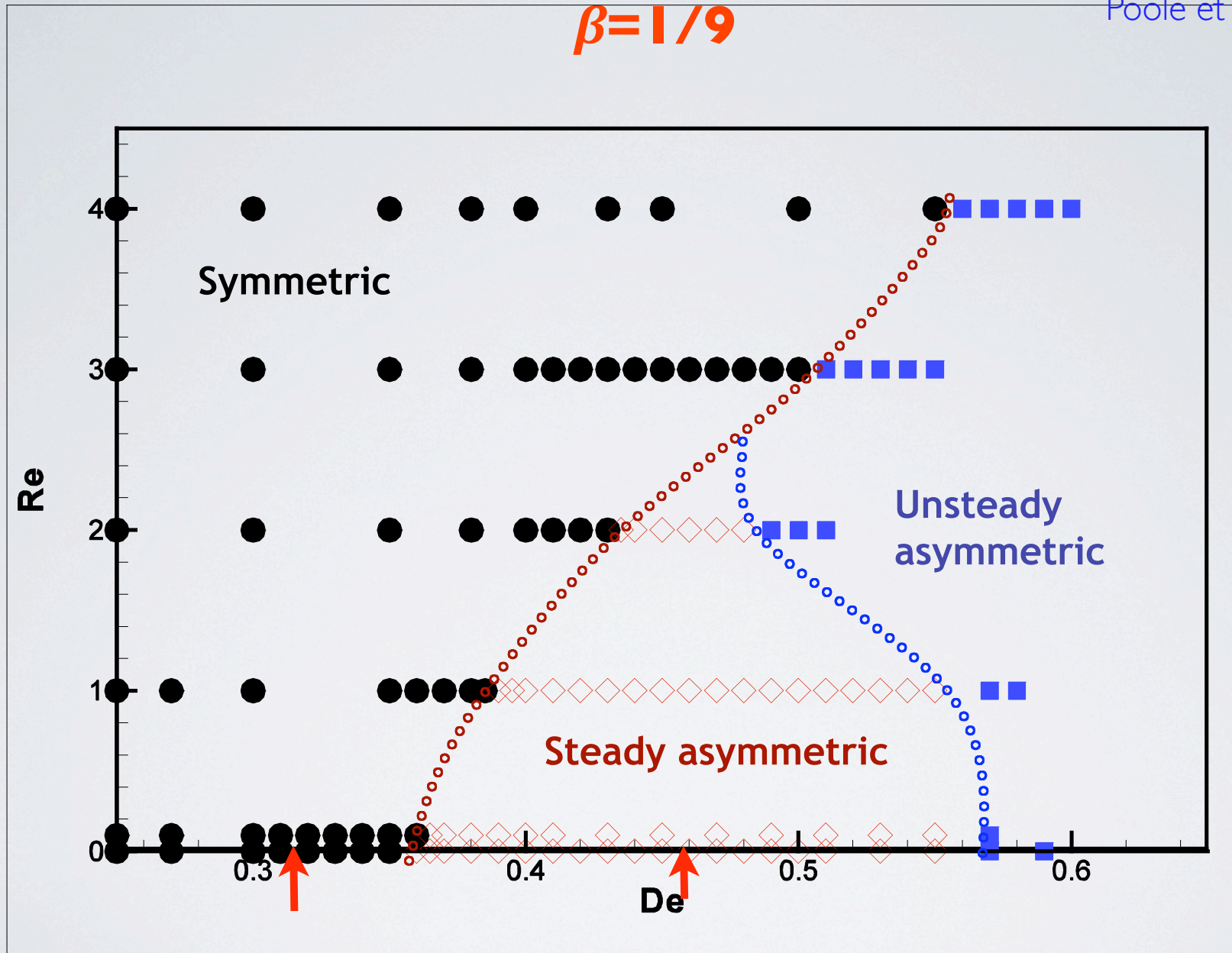


Transitions in some stagnation viscoelastic flows at $Re=0$
Flow Instabilities and turbulence in viscoelastic fluids

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2D CROSS SLOT: OLDROYD-B — STABILITY MAP

Poole et al., SoR 2007



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2D CROSS SLOT: SPTT — EFFECT OF EPSILON

Poole et al., SoR 2007

$$\beta = 1/9$$

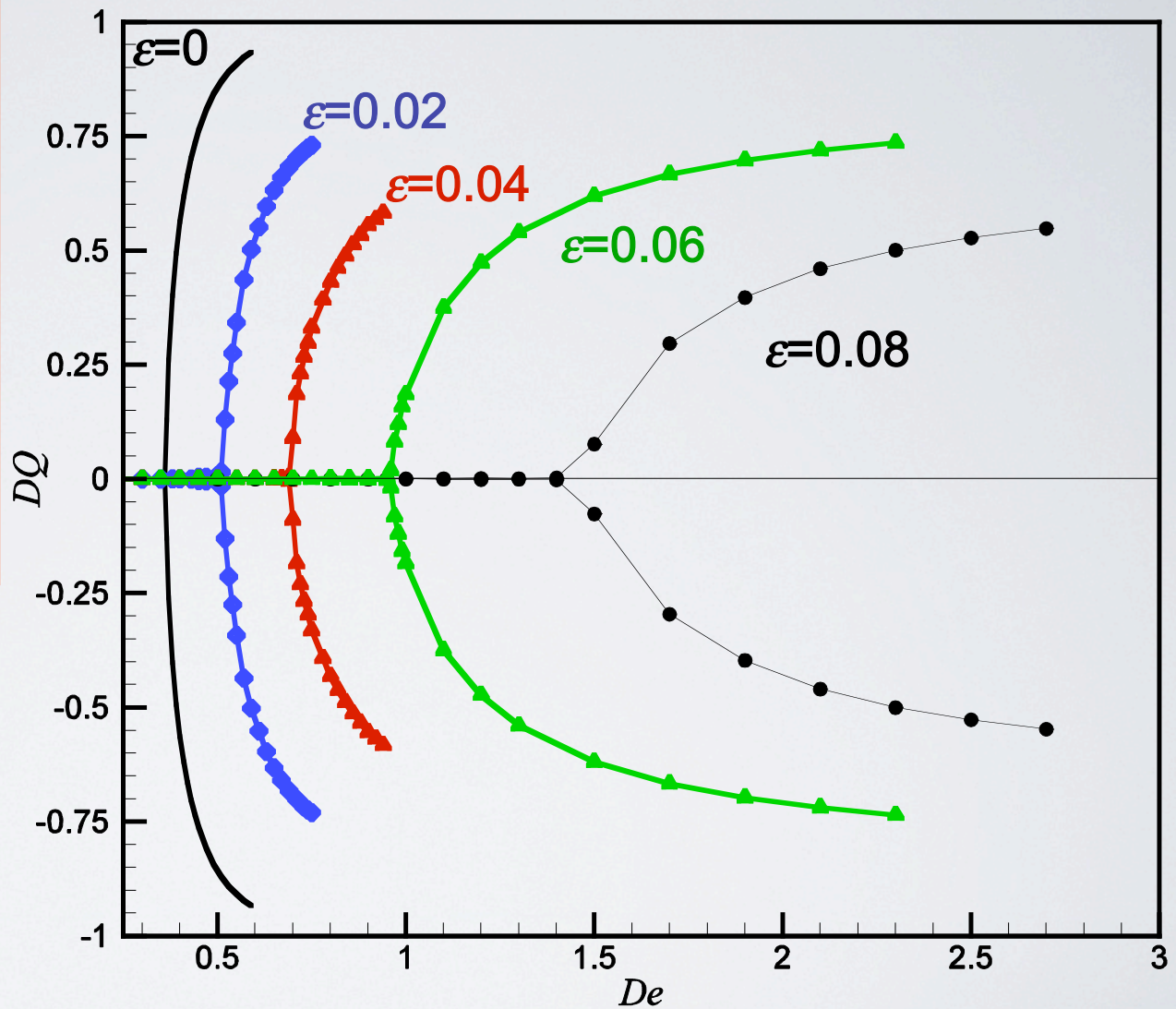
Increasing ε

Increases De_{CR}

Decreases degree of asymmetry ($\varepsilon < 0.04$)

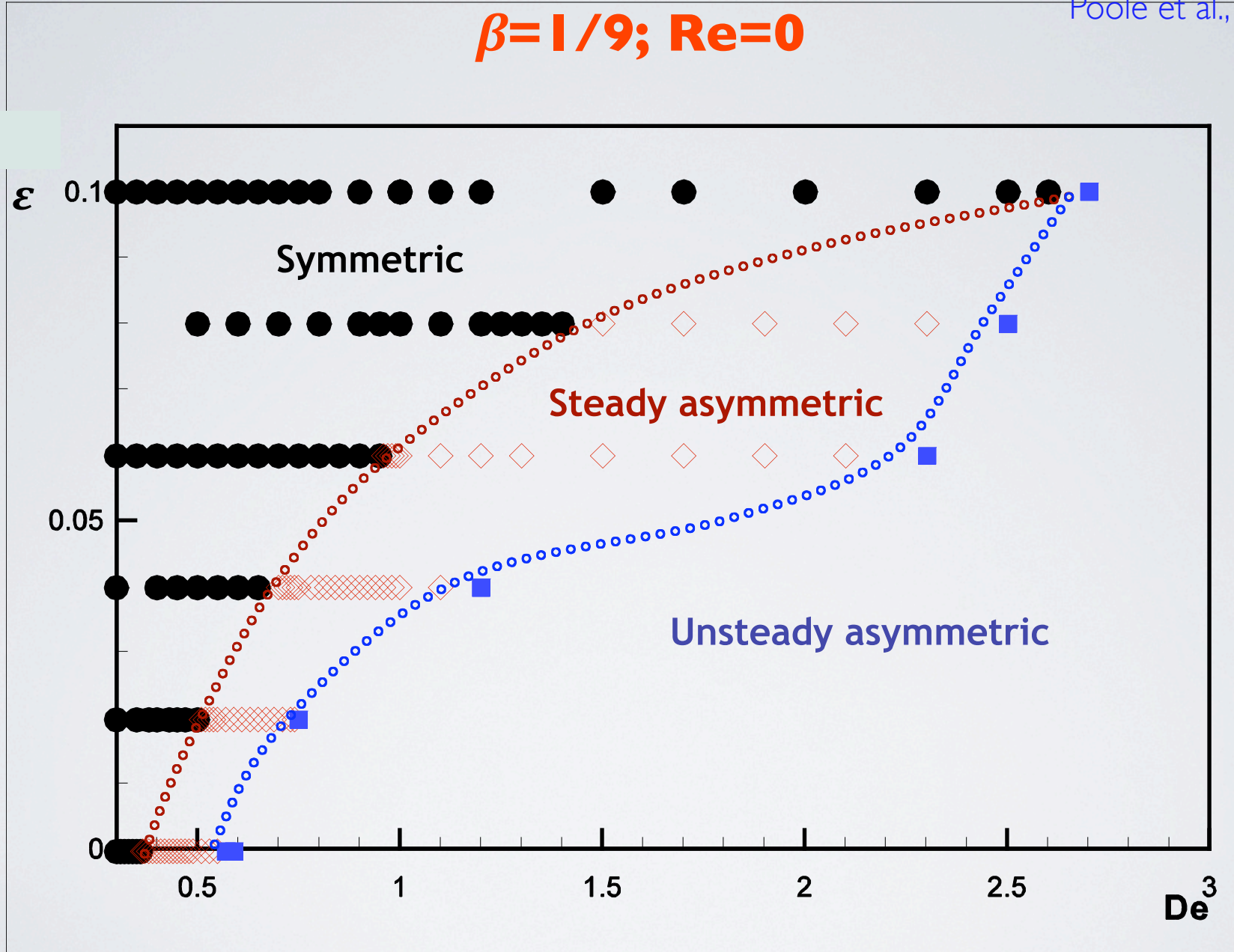
Increases degree of asymmetry and extension in De ($\varepsilon > 0.04$)

Asymmetric stable flow disappears for $\varepsilon > 0.08$



2D CROSS SLOT: SPTT — STABILITY MAP

Poole et al., SoR 2007



Transitions in some stagnation viscoelastic flows at $Re=0$
Flow Instabilities and turbulence in viscoelastic fluids

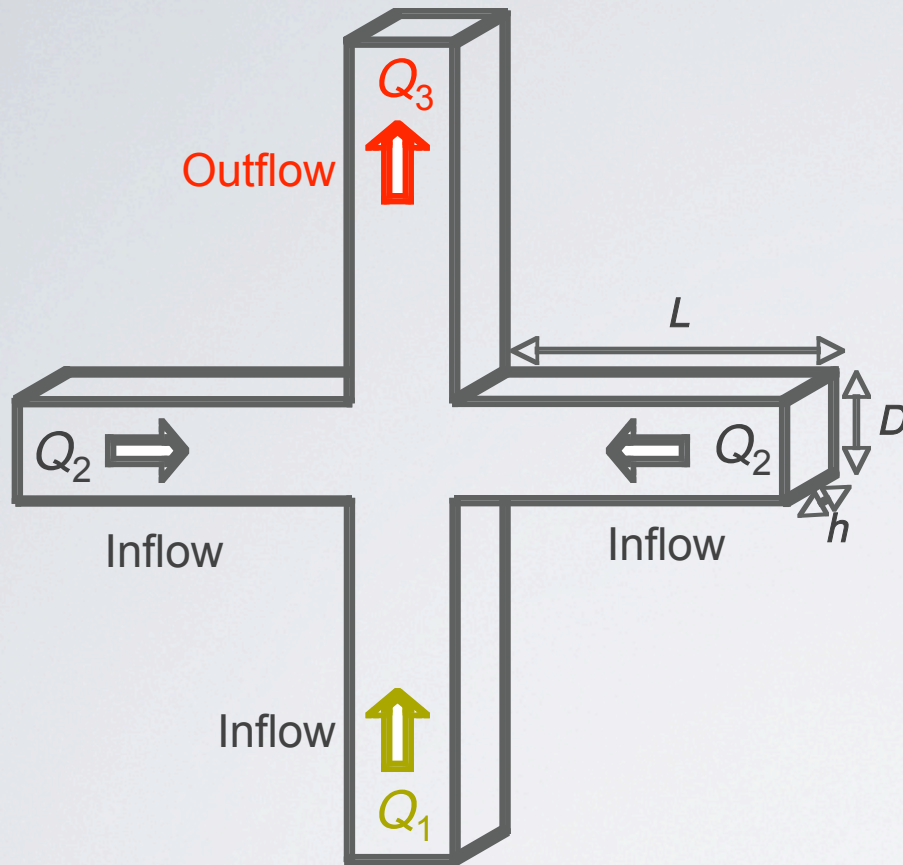
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FLOW FOCUSING

(extensional flow “without” shear)

FLOW FOCUSING

Oliveira et al. JNNFM 160 (2009) 31-39



**Cross-slot with
3 inlets and 1 outlet**

Operational Variables

$$Q_1, Q_2$$

$$Q_3 = 2 \times Q_2 + Q_1$$

Dimensionless Variables

$$FR = \frac{Q_2}{Q_1}$$

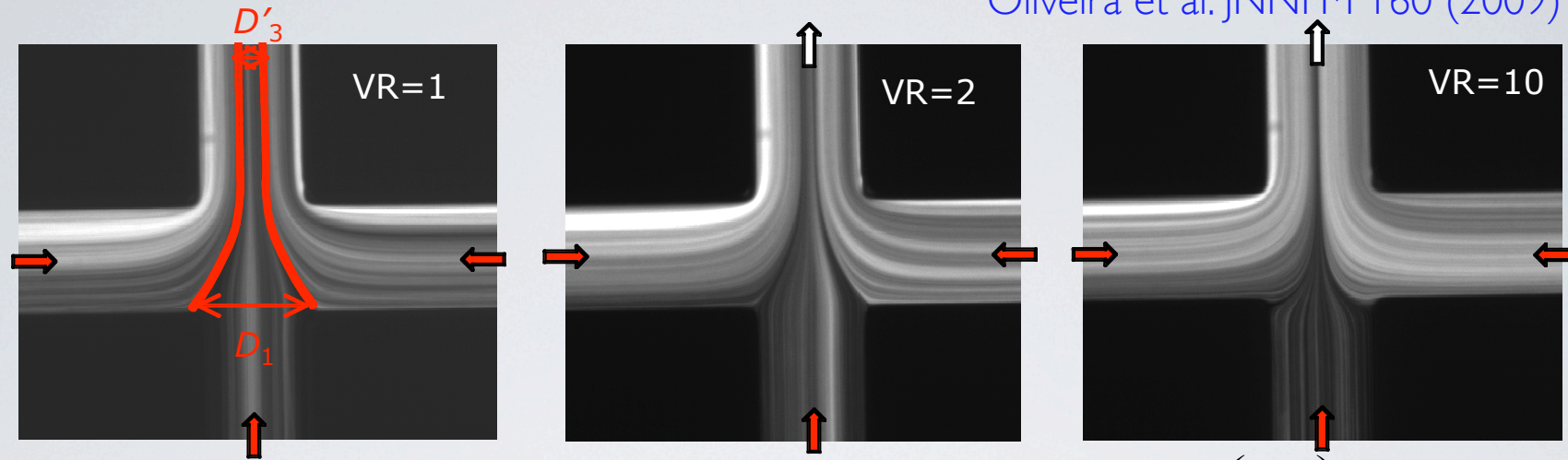
$$VR = \frac{U_2}{U_1} \quad (= FR)$$

$$\left. \begin{aligned} Re &= \frac{\rho U_2 D}{\eta_0} \\ De &= \frac{\lambda U_2}{D} \end{aligned} \right\} El = \frac{De}{Re}$$

All dimensions kept constant
in experiments and calculations

FLOW FOCUSING: NEWTONIAN

Oliveira et al. JNNFM 160 (2009) 31-39

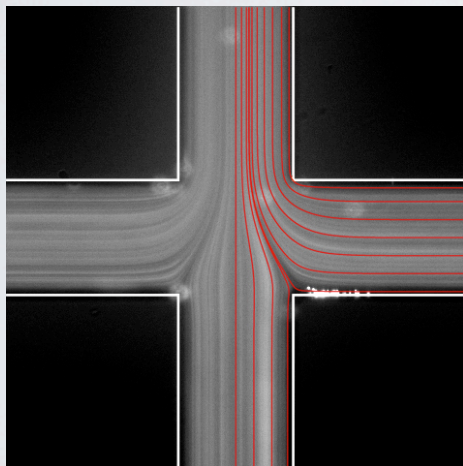


Separation streamlines: nearly hyperbolic shape

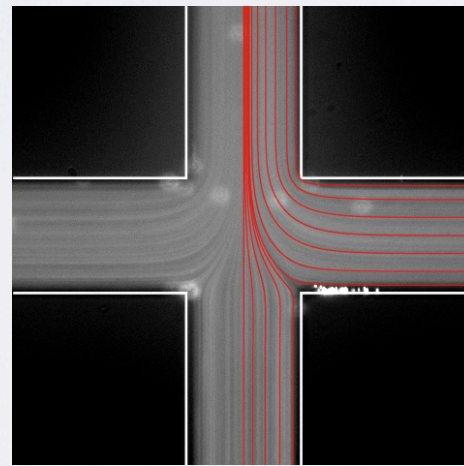
$$\varepsilon_H = \ln\left(\frac{D_1}{D_3}\right) = \ln\left[\frac{3}{2}(1+2VR)\right]$$

$Q_1 = 0.01$ ml/h

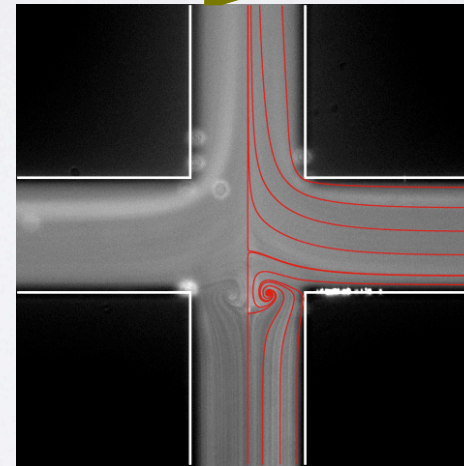
increasing Q_2



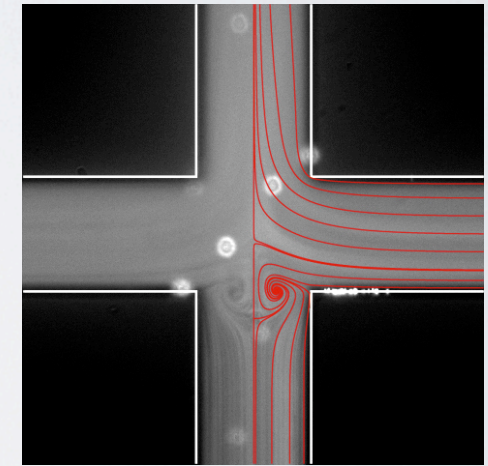
$Q_2 = 0.3$ ml/h
 $VR = 1, Re_3 = 2.8$



$Q_2 = 0.9$ ml/h
 $VR = 3, Re_3 = 6.5$



$Q_2 = 15$ ml/h
 $VR = 50, Re_3 = 94.2$



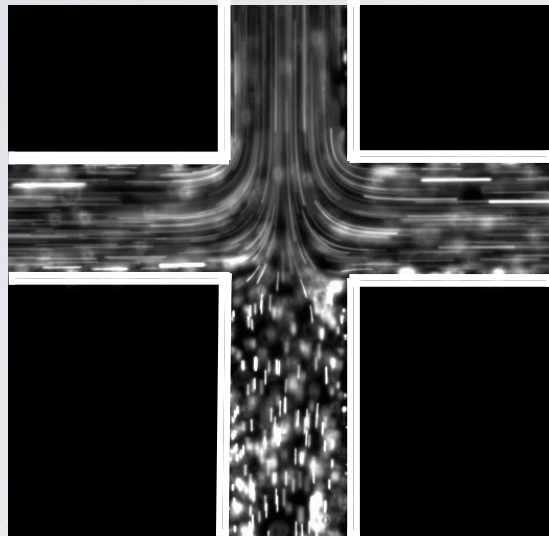
$Q_2 = 18$ ml/h
 $VR = 60, Re_3 = 112.8$

$Q_1 = 0.01$ ml/h

increasing Q_2

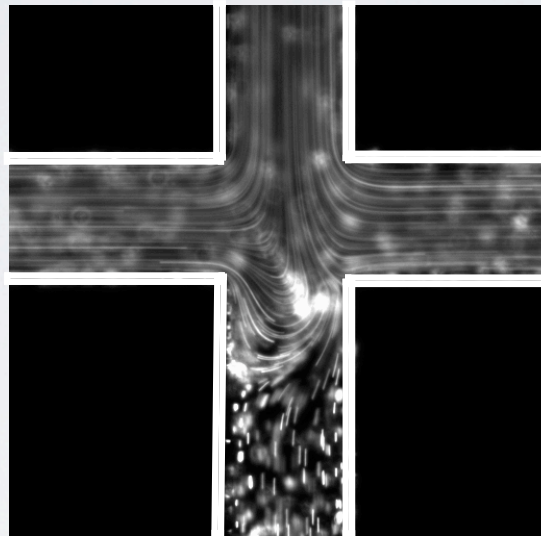


Viscoelastic



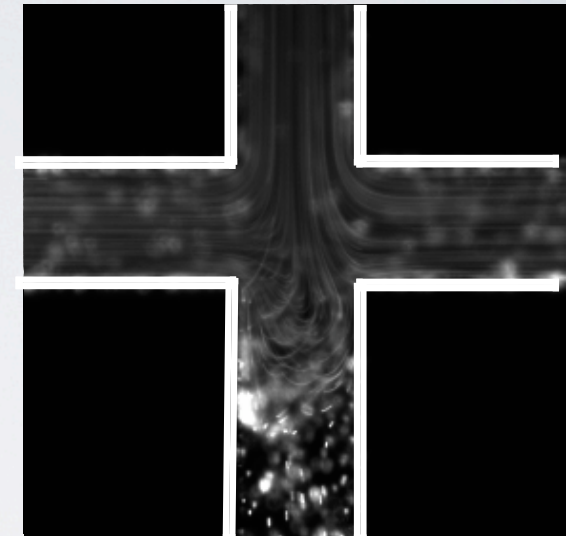
$Q_2 = 0.05$ ml/h, $VR = 5$
 $Re = 0.23$, $De = 0.38$

Symmetric



$Q_2 = 0.2$ ml/h, $VR = 20$
 $Re = 0.87$, $De = 1.41$

Steady Asymmetric



$Q_2 = 0.5$ ml/h, $VR = 50$
 $Re = 2.15$, $De = 3.479$

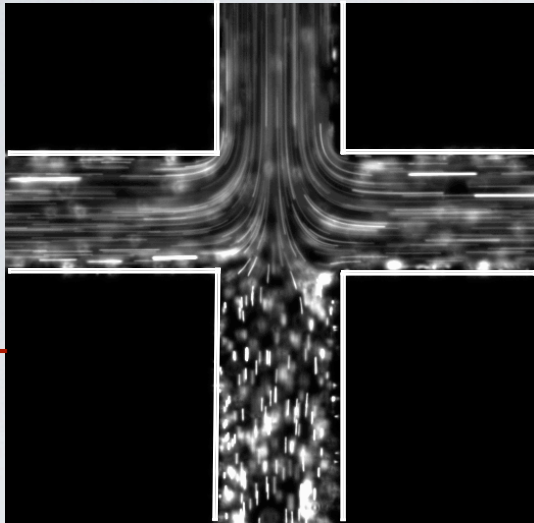
Unsteady 3D

FLOW FOCUSING: PAA 125

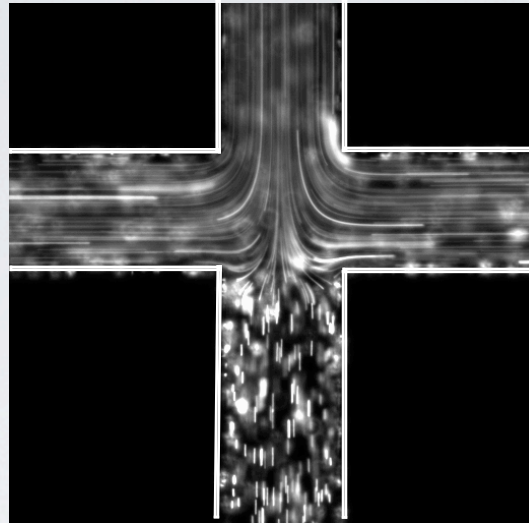
$Q_1 = 0.01$ ml/h

Oliveira et al. JNNFM 160 (2009) 31-39

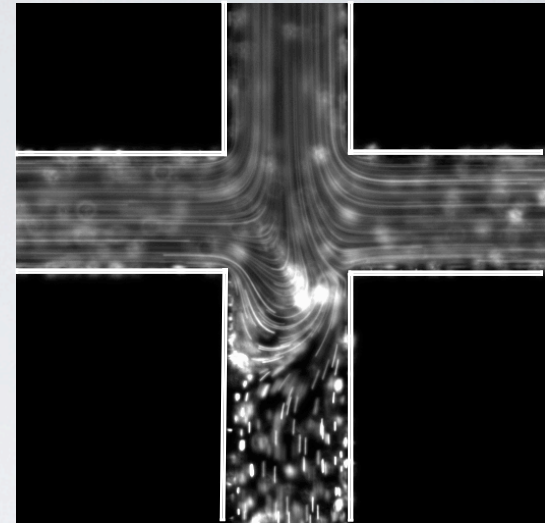
Viscoelastic
Experimental



$Q_2 = 0.05$ ml/h, $VR = 5$
 $Re = 0.23$, $De = 0.38$

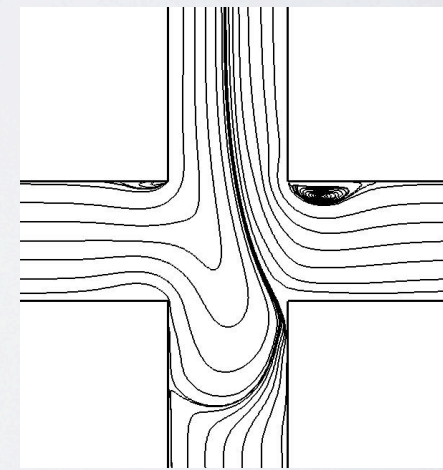
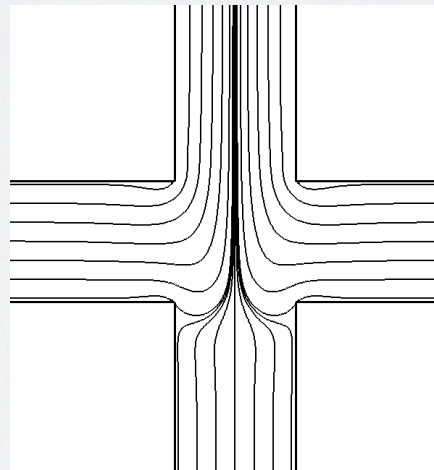
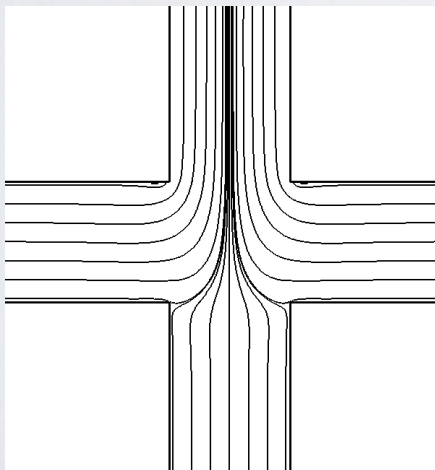


$Q_2 = 0.1$ ml/h, $VR = 10$
 $Re = 0.45$, $De = 0.723$



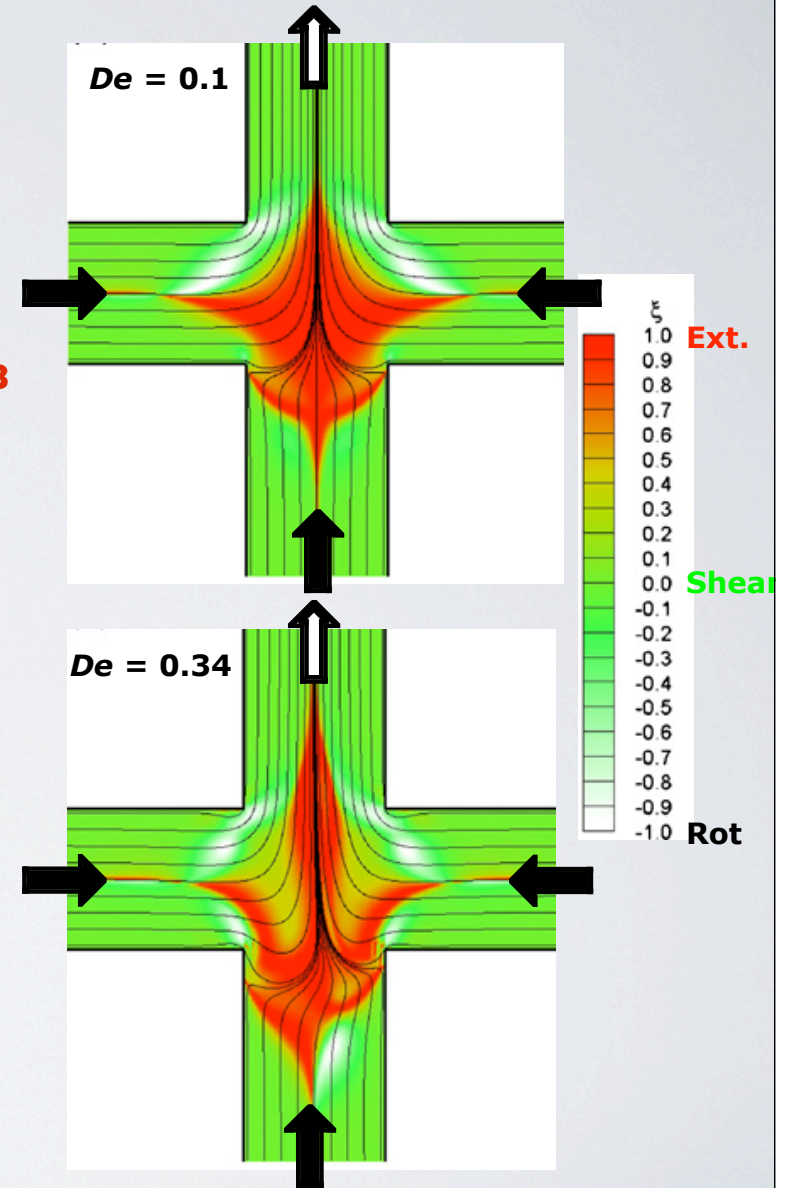
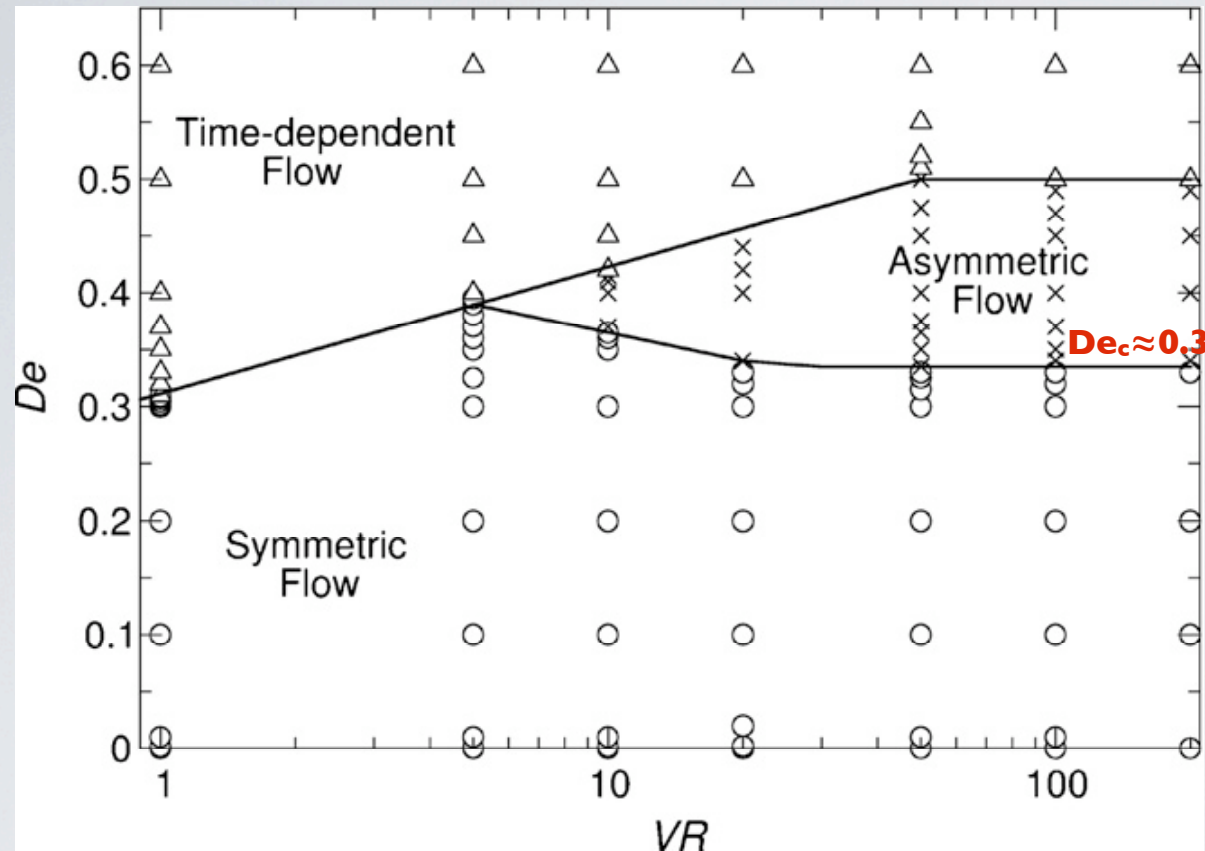
$Q_2 = 0.2$ ml/h, $VR = 20$
 $Re = 0.87$, $De = 1.41$

UCM
2D Calculations



FLOW FOCUSING: VISCOELASTIC

Oliveira et al. JNNFM 160 (2009) 31-39



Astarita, JNNFM 6 (1979) 69
 Thompson et al., JNNFM 86 (1999) 375
 Mompean et al., JNNFM 111 (2003) 151

$$\xi = \frac{1 - R}{1 + R}$$

$$R = \frac{tr \tilde{\mathbf{W}}^2}{tr \mathbf{D}^2}$$

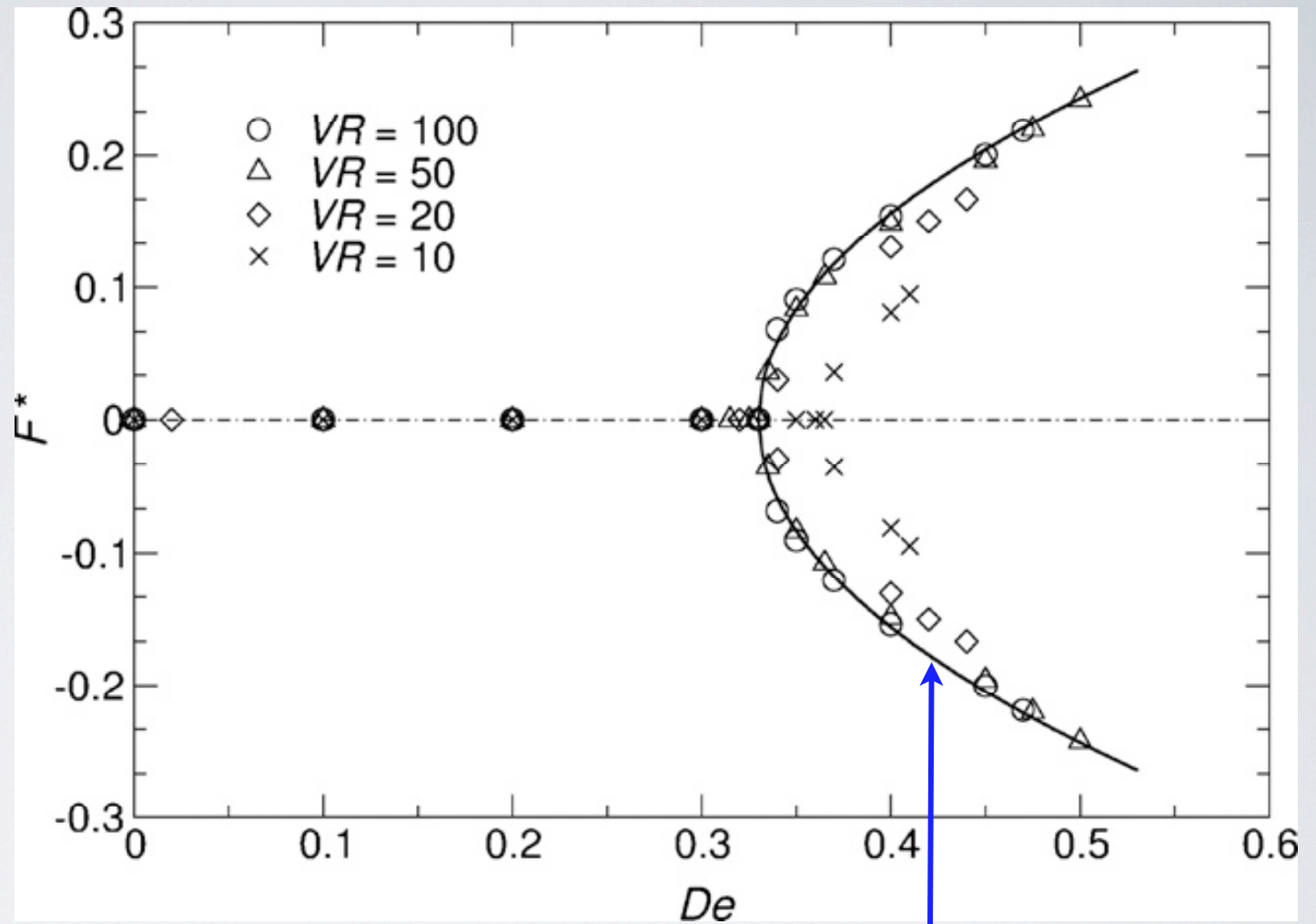
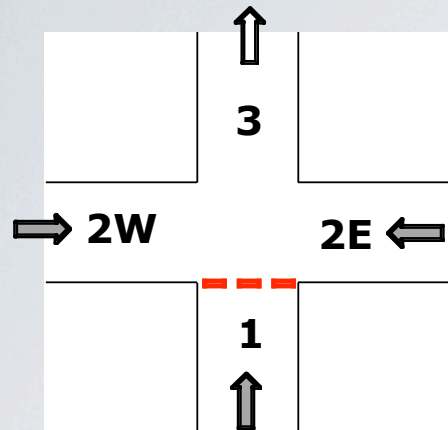
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FLOW FOCUSING: EFFECT OF VR

Oliveira et al. JNNFM 160 (2009) 31-39

$$F^* = \frac{F_W - F_E}{F_3}$$



Bistable flow

High VR:

constant De_c

evolution independent of VR

supercritical pitchfork bifurcation

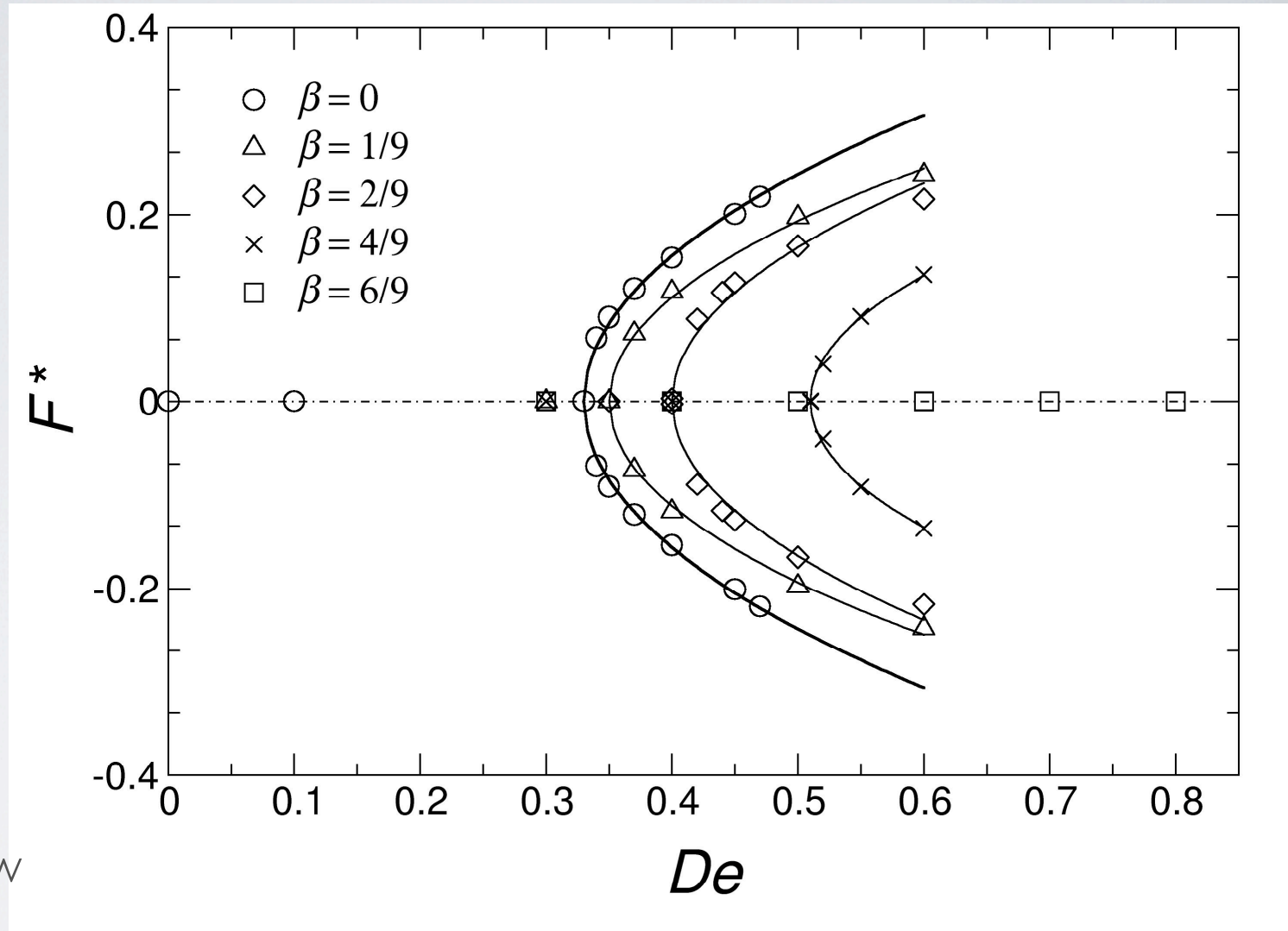
$$F^* = 0.59\sqrt{De - 0.33}$$

FLOW FOCUSING: EFFECT OF β

Oliveira et al. JNNFM 160 (2009) 31-39

$$\beta = \frac{\eta_s}{\eta_s + \eta_p}$$

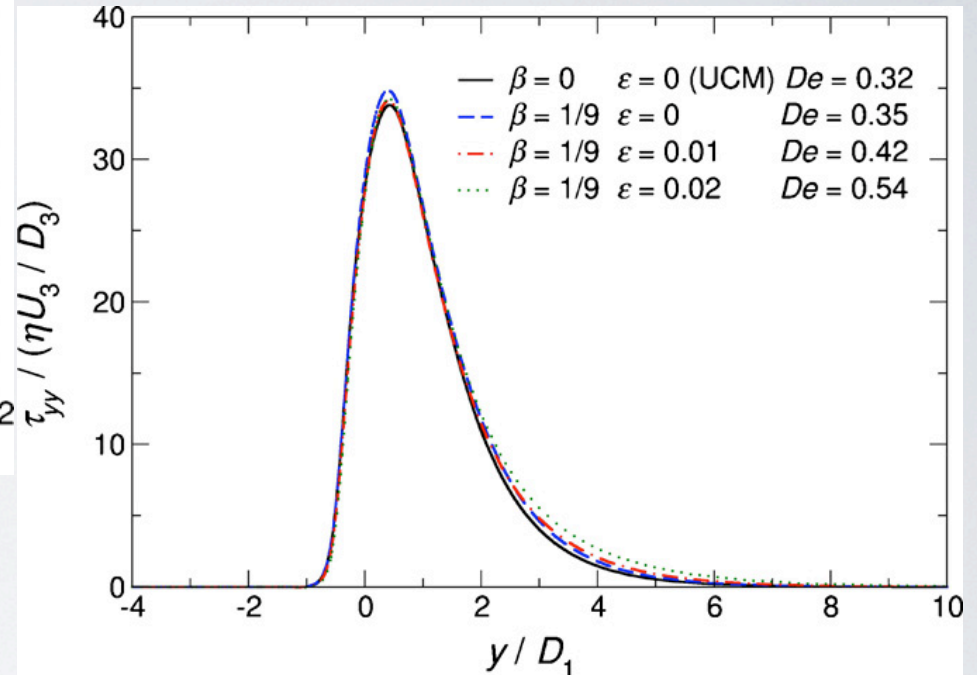
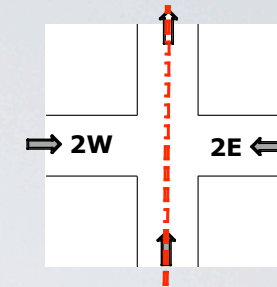
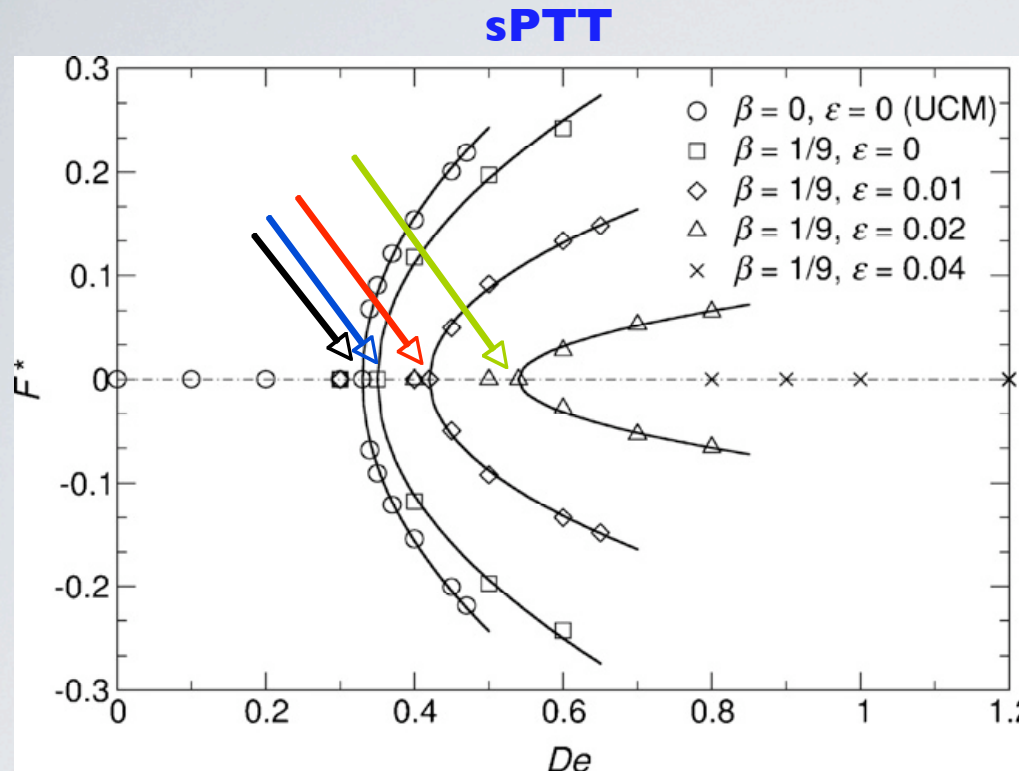
↓
Oldroyd-B



β stabilizes the flow
increases De_c

$\beta \geq 6/9$, no steady asymmetry

FLOW FOCUSING: EFFECT OF ε

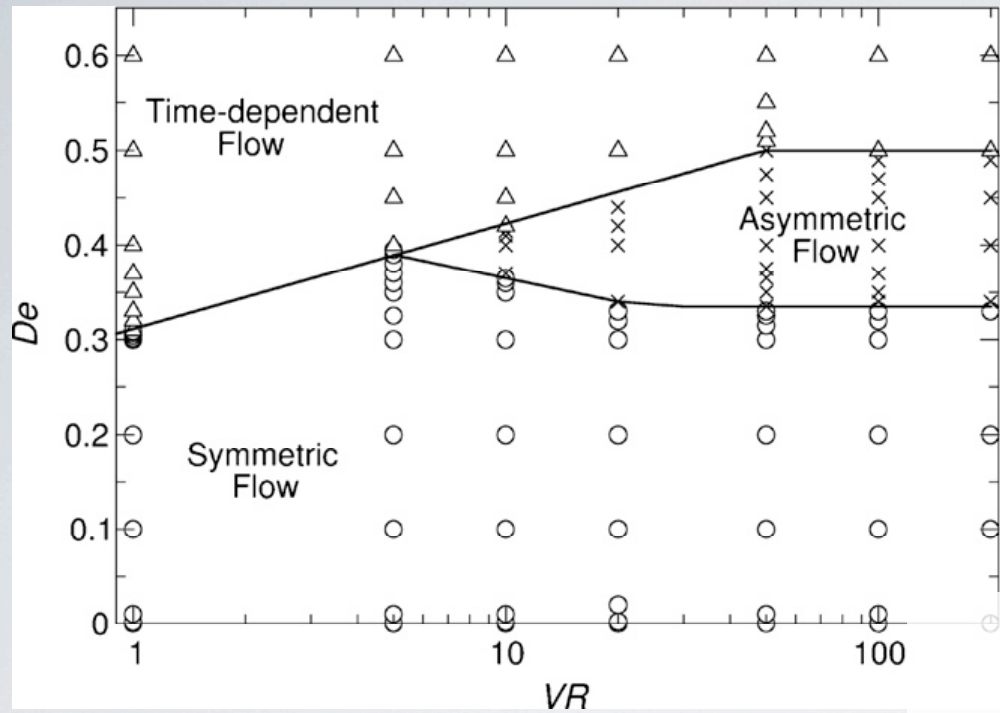


ε stabilizes the flow
 increases De_c
 decreases degree of asymmetry
 $\varepsilon \geq 0.04$ steady asymmetry disappears
(Transition directly to unsteady flow)

Similar levels of normal stresses
 achieved near critical conditions
 Extensional properties decisive
 for onset of flow asymmetry

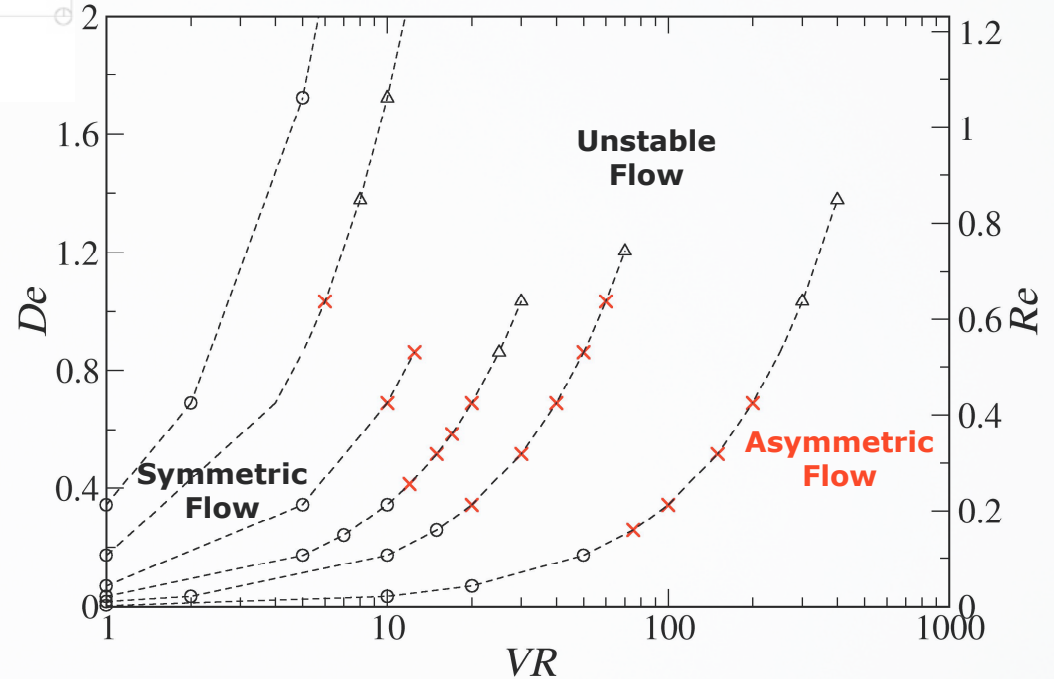
FLOW FOCUSING: NUMERICAL VERSUS EXPERIMENTS (PAA 125)

Oliveira et al. JNNFM 160 (2009) 31-39



Numerical
UCM, 2D, $Re=0$

Experimental
PAA 125 + NaCl



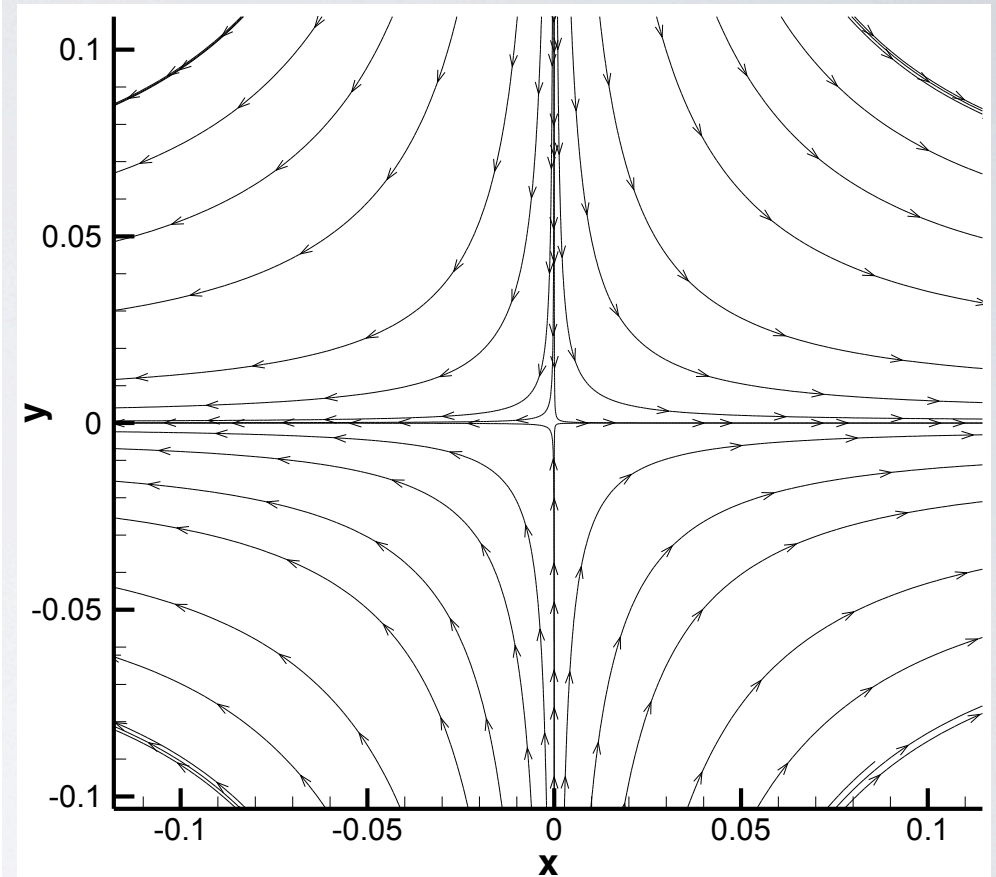
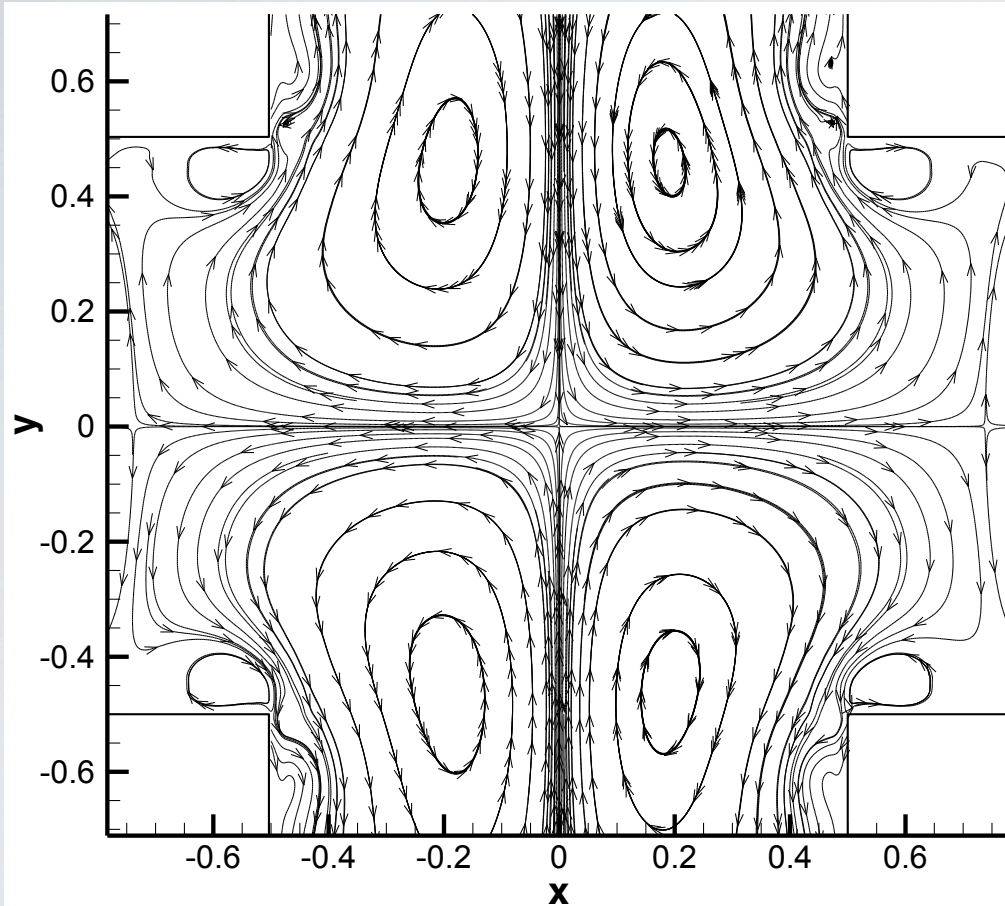
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STAGNATION FLOW
Symmetry & asymmetry
**Some observations from numerics on cross
flow**

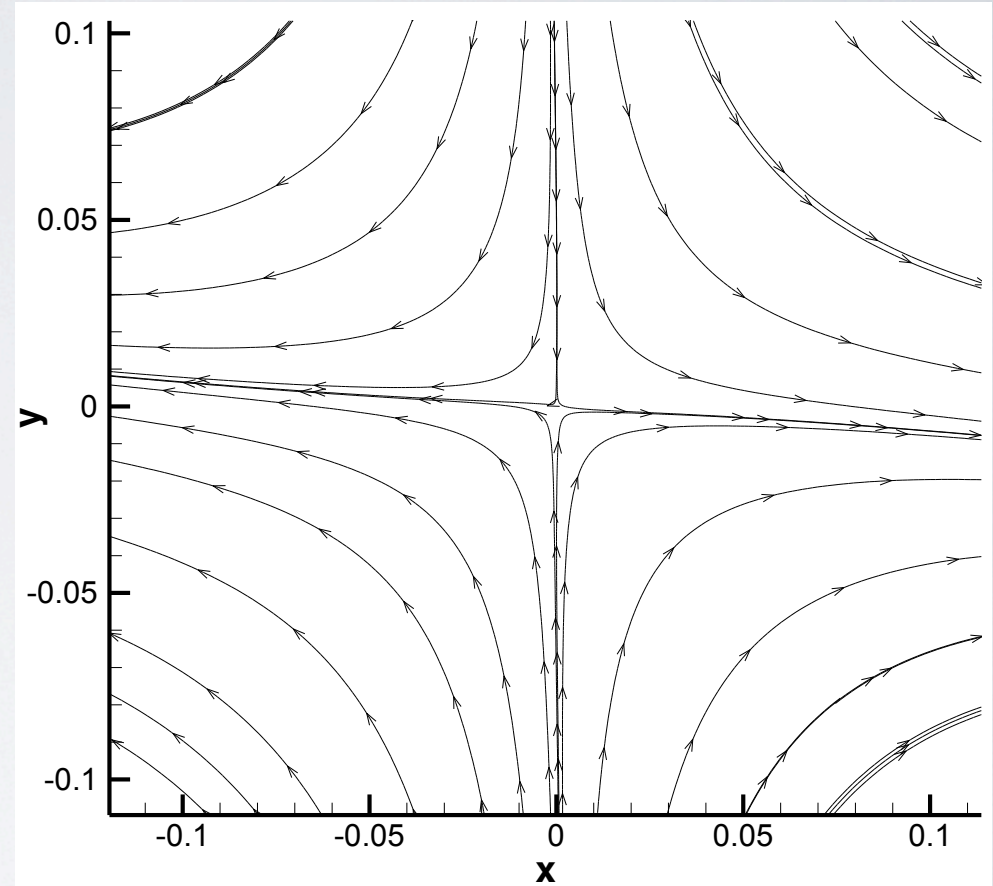
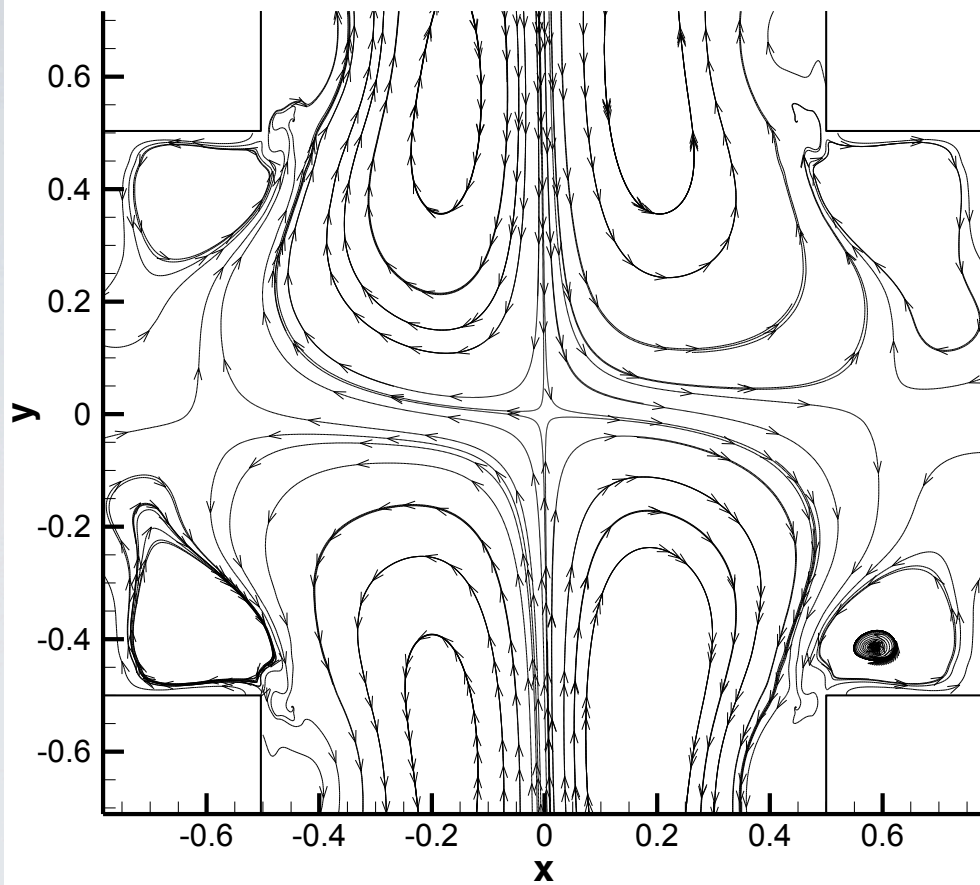
DIFFERENCE OF TWO SYMMETRIC FLOWS FAR FROM TRANSITION

$$De = 0.20 - De = 0.19$$



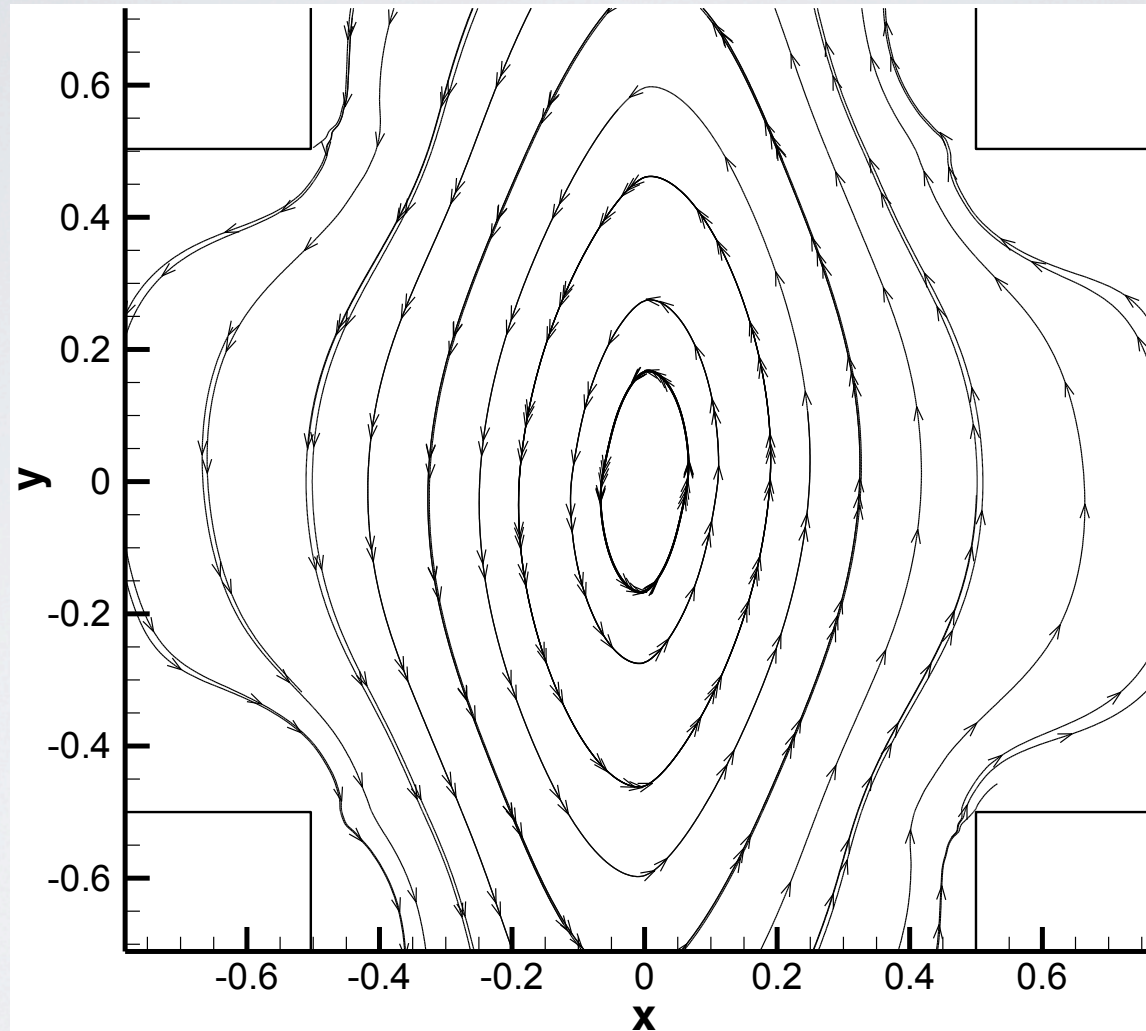
DIFFERENCE OF TWO SYMMETRIC FLOWS CLOSE TO TRANSITION

$$De = 0.308 - De = 0.307$$



CRITICAL FLOW - SYMMETRIC FLOW

$$De = 0.310 - De = 0.309$$

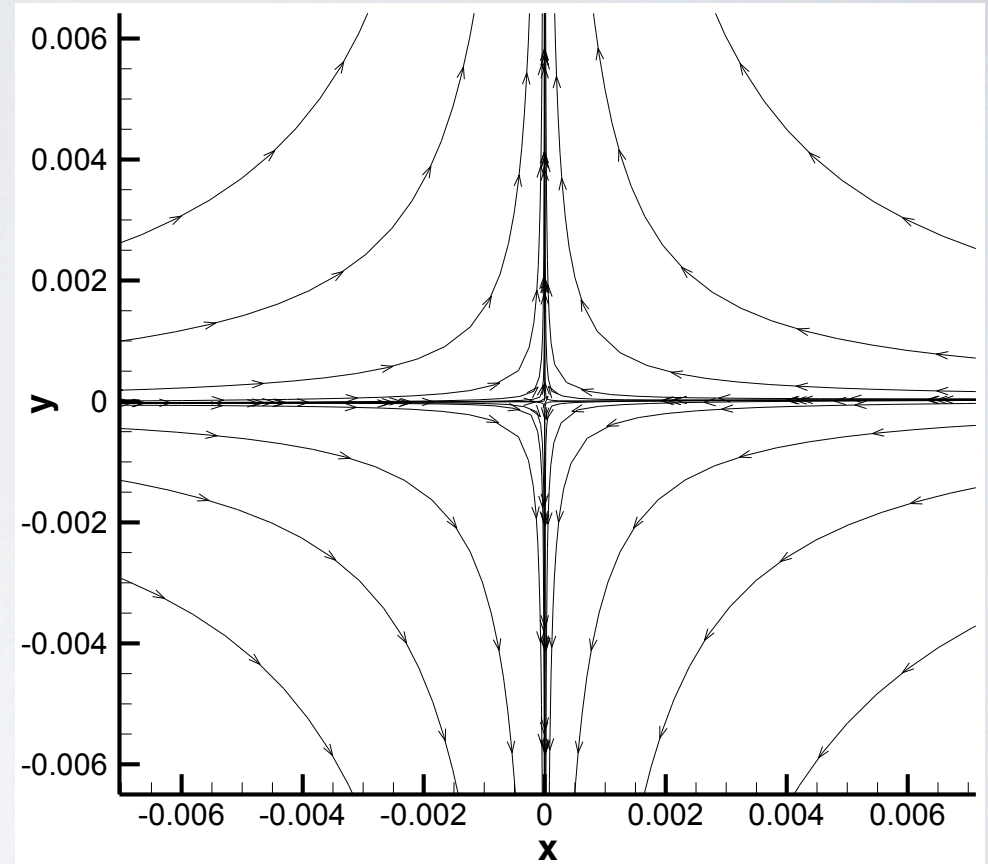
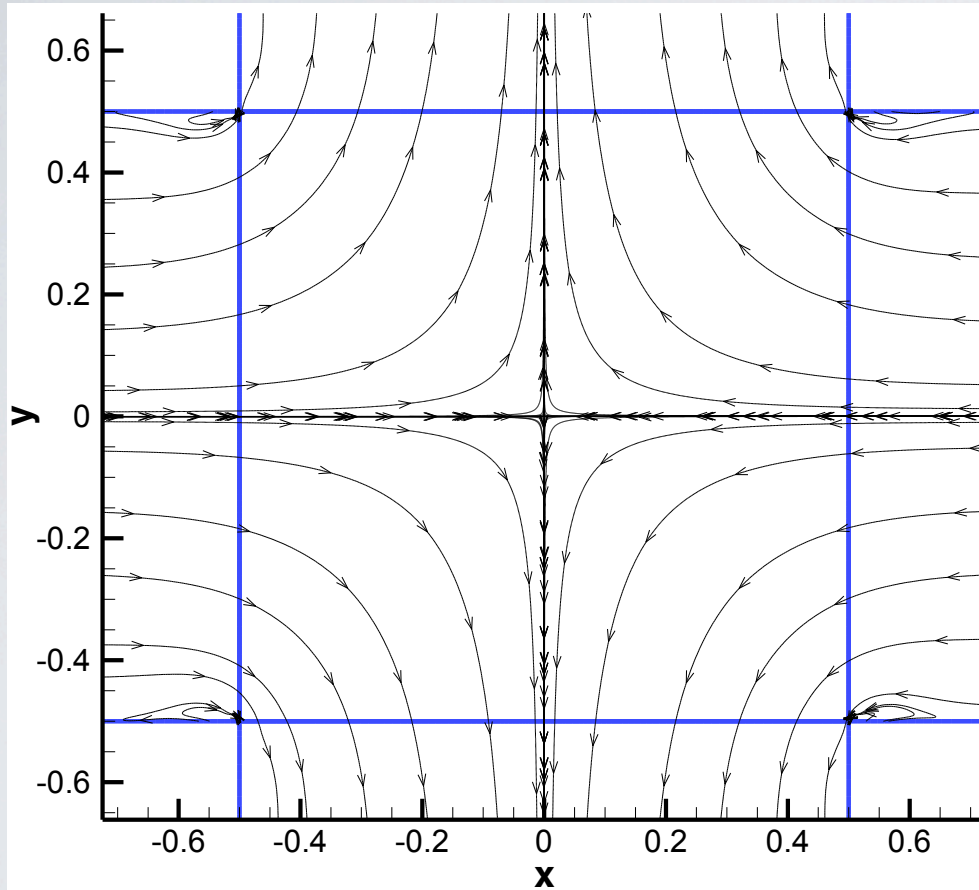


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STREAMLINES FOR SYMMETRIC FLOW

$$De = 0.309$$

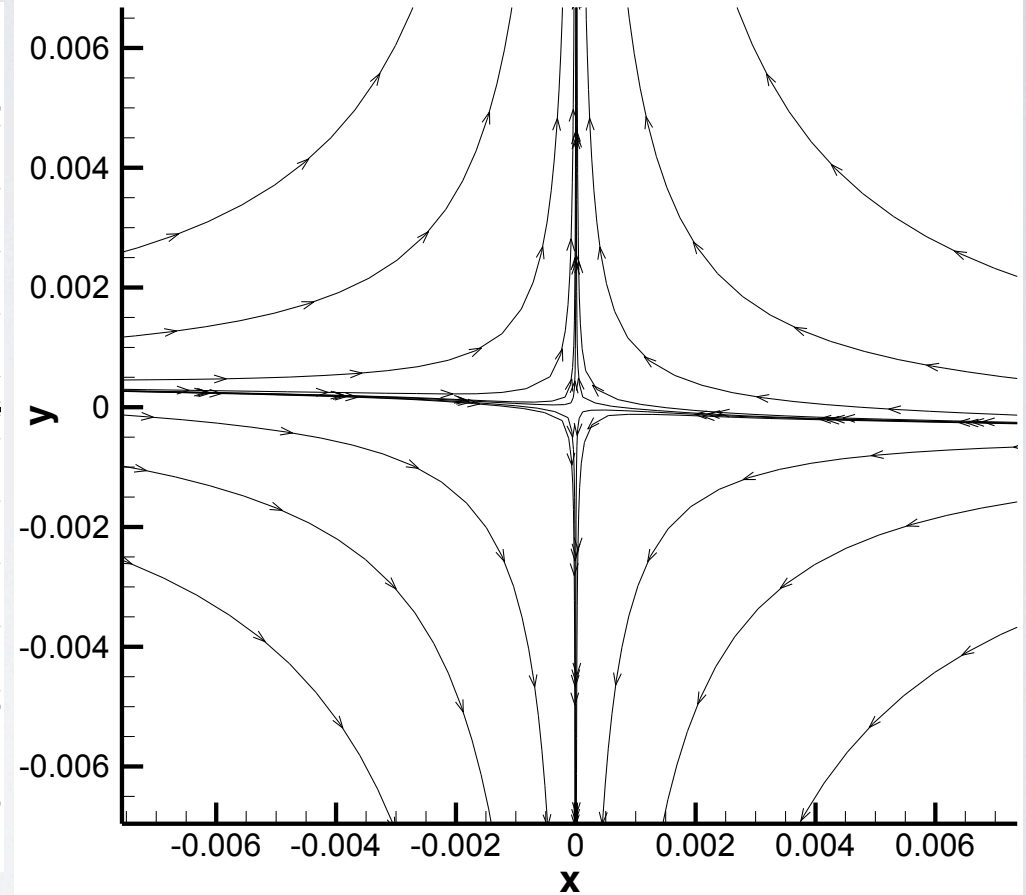
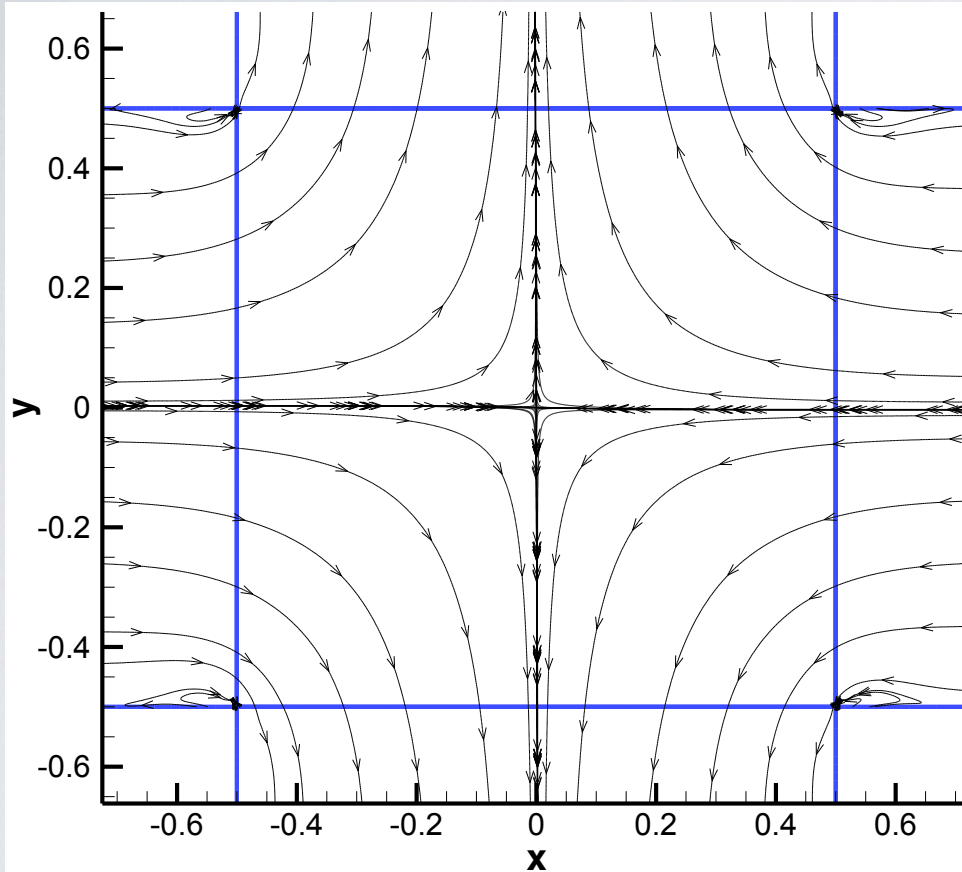


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STREAMLINES FOR CRITICAL FLOW

$$De = 0.310$$



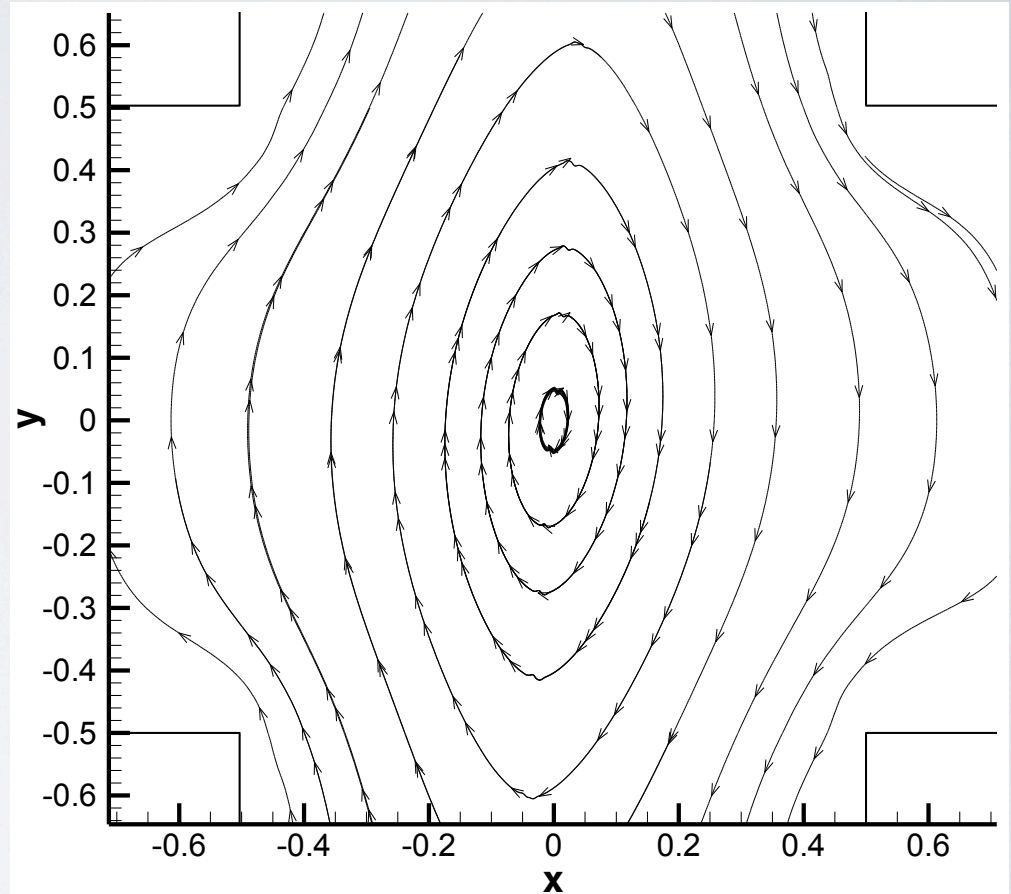
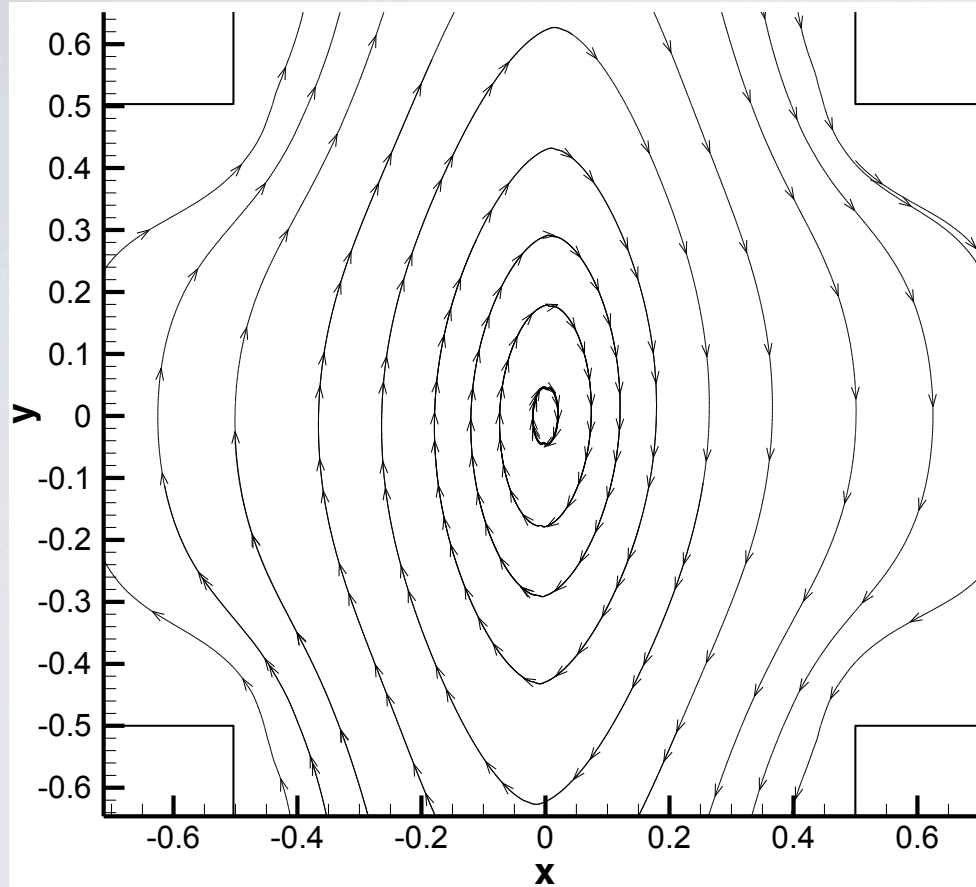
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SMALL DIFFERENCE BETWEEN TWO ASYMMETRIC FLOWS

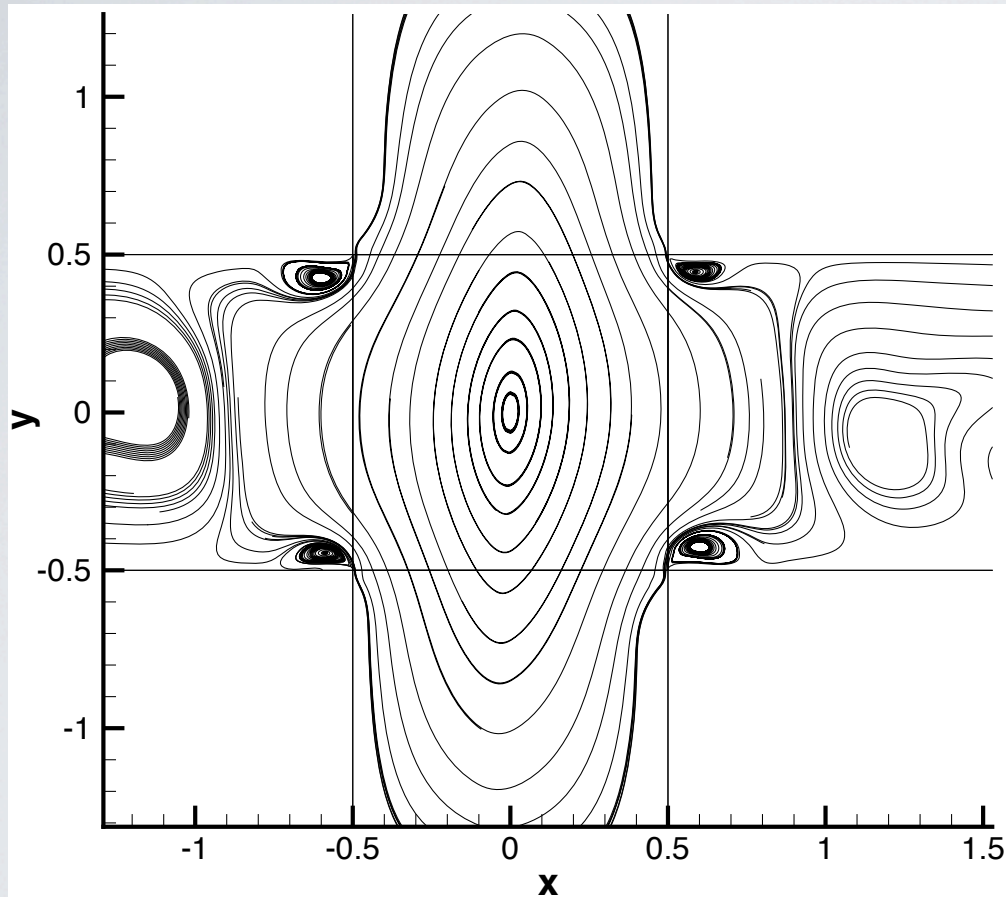
$$De = 0.312 - De = 0.311$$

$$De = 0.315 - De = 0.314$$

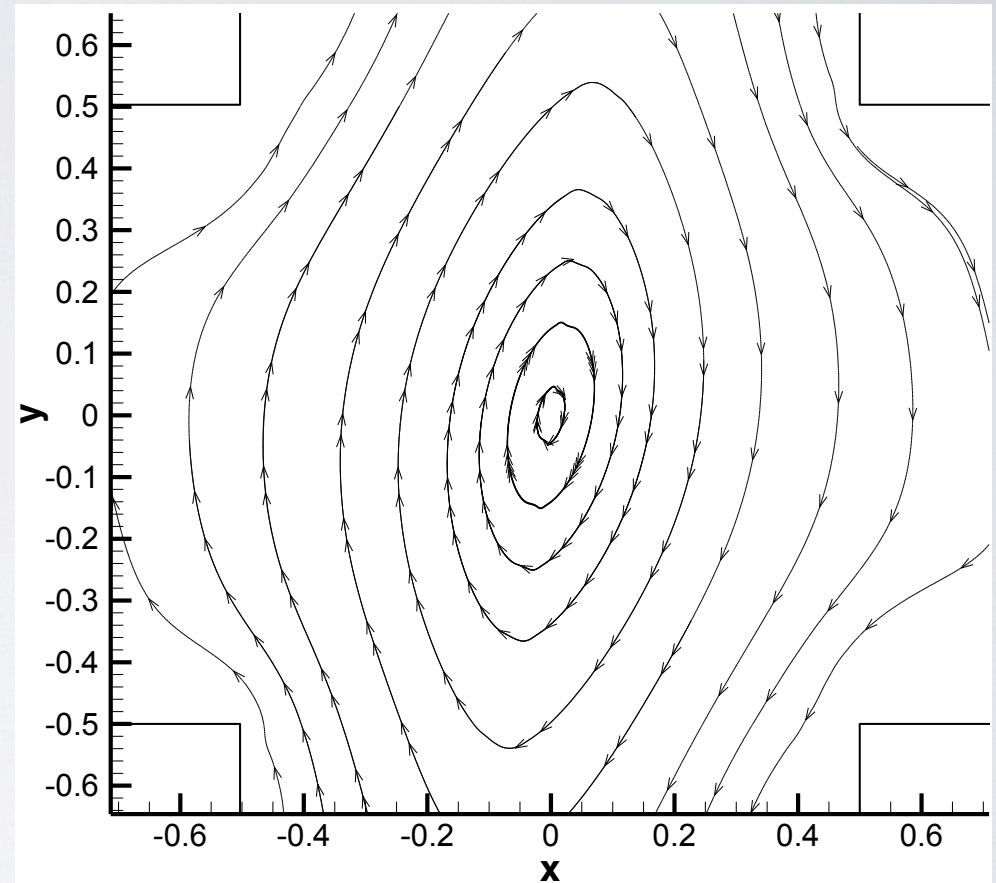


LARGER DIFFERENCES

Asymmetric- critical
 $De = 0.32 - De = 0.31$



Asymmetric- asymmetric
 $De = 0.34 - De = 0.32$



STAGNATION + VORTEX FLOW

An analytical solution

PROBLEM FORMULATION: UCM

Stagnation flow

$$u_{sta} = ax$$

$$v_{sta} = -ay$$

“Vortex” flow

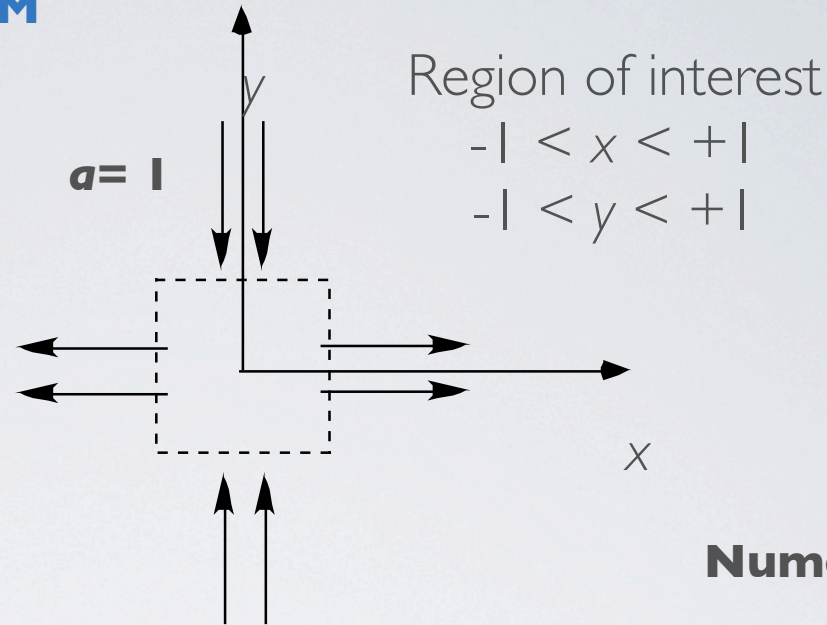
$$u_{vor} = b_u y$$

$$v_{vor} = b_v x$$

Stagnation + “vortex” flow

$$u = ax + b_u y$$

$$v = -ay + b_v x$$



$$\tau_{xx} + De \left[u \frac{\partial \tau_{xx}}{\partial x} + v \frac{\partial \tau_{xx}}{\partial y} - 2 \left(\tau_{xx} \frac{\partial u}{\partial x} + \tau_{xy} \frac{\partial v}{\partial x} \right) \right] = 2 \frac{\partial u}{\partial x}$$

$$\tau_{xy} + De \left[u \frac{\partial \tau_{xy}}{\partial x} + v \frac{\partial \tau_{xy}}{\partial y} - \left(\tau_{xx} \frac{\partial u}{\partial y} + \tau_{yy} \frac{\partial v}{\partial x} \right) \right] = \frac{\partial u}{\partial y} + \frac{\partial v}{\partial x}$$

$$\tau_{yy} + De \left[u \frac{\partial \tau_{yy}}{\partial x} + v \frac{\partial \tau_{yy}}{\partial y} - 2 \left(\tau_{xy} \frac{\partial u}{\partial y} + \tau_{yy} \frac{\partial v}{\partial y} \right) \right] = 2 \frac{\partial v}{\partial y}$$

GENERAL SOLUTION & CONSTANT SOLUTION

$$\tau_{xx} + De \left[(b_u y + ax) \frac{\partial \tau_{xx}}{\partial x} + (b_v x - ay) \frac{\partial \tau_{xx}}{\partial y} - 2(a\tau_{xx} + b_v \tau_{xy}) \right] = 2a$$

$$\tau_{xy} + De \left[(b_v x - ay) \frac{\partial \tau_{xy}}{\partial x} + (b_u y + ax) \frac{\partial \tau_{xy}}{\partial y} - (b_u \tau_{xx} + b_v \tau_{yy}) \right] = b_u + b_v$$

$$\tau_{yy} + De \left[(b_v x - ay) \frac{\partial \tau_{yy}}{\partial x} + (b_u y + ax) \frac{\partial \tau_{yy}}{\partial y} - 2(b_u \tau_{xy} - a\tau_{yy}) \right] = -2a$$

$$\tau_{ij} = (\tau_{ij})|_{const} + (\tau_{ij})|_{homogeneous}$$

$$\frac{\partial \tau_{ij}}{\partial x_k} = 0 \Rightarrow \tau_{xx} = -\frac{2[a + 2a^2 De + b_v(b_u + b_v)De]}{-1 + 4(a^2 + b_u b_v)De^2}$$

$$\tau_{xy} = -\frac{b_u + b_v + 2a(b_u - b_v)De}{-1 + 4(a^2 + b_u b_v)De^2} \quad \tau_{yy} = \frac{2[a - 2a^2 De - b_u(b_u + b_v)De]}{-1 + 4(a^2 + b_u b_v)De^2}$$

This solution absorbs the constants on the rhs of constitutive equation

HOMOGENEOUS SOLUTION (I)

$$\tau_{xx} + De \left[(b_u y + ax) \frac{\partial \tau_{xx}}{\partial x} + (b_v x - ay) \frac{\partial \tau_{xx}}{\partial y} - 2(a\tau_{xx} + b_v \tau_{xy}) \right] = 0$$

$$\tau_{xy} + De \left[(b_v x - ay) \frac{\partial \tau_{xy}}{\partial x} + (b_u y + ax) \frac{\partial \tau_{xy}}{\partial y} - (b_u \tau_{xx} + b_v \tau_{yy}) \right] = 0$$

$$\tau_{yy} + De \left[(b_v x - ay) \frac{\partial \tau_{yy}}{\partial x} + (b_u y + ax) \frac{\partial \tau_{yy}}{\partial y} - 2(b_u \tau_{xy} - a\tau_{yy}) \right] = 0$$

Solution hypothesis (I): $\tau_{ij}(x, y) = \tau_{ij}(\phi)$ with $\phi = kx + Ty$

$$m\phi De \sqrt{a^2 + b_u b_v} \frac{d\tau_{xx}}{d\phi} = (-1 + 2aDe)\tau_{xx} + 2b_v De \tau_{xy}$$

$$m\phi De \sqrt{a^2 + b_u b_v} \frac{d\tau_{xy}}{d\phi} = b_u De \tau_{xx} + b_v De \tau_{yy}$$

$$m\phi De \sqrt{a^2 + b_u b_v} \frac{d\tau_{yy}}{d\phi} = -(1 + 2aDe)\tau_{yy} + 2b_u De \tau_{xy}$$

$$k = \frac{Tb_v}{-a \pm \sqrt{a^2 + b_u b_v}}$$

$$m = \pm 1$$

HOMOGENEOUS SOLUTION (2)

Solution hypothesis (2): $\tau_{ij}(\phi) = \alpha_{ij}\phi^q$

as in stagnation flow ^{1,2}

¹ Renardy JNNFM 138 (2006) 204-205

² Becherer, Morozov, van Saarloos JNNFM 153 (2008) 183-190

$$\left[\left(-1 + 2aDe - mqDe\sqrt{a^2 + b_u b_v} \right) \alpha_{xx} + 2b_v De \alpha_{xy} \right] \phi^q = 0$$

$$\left[b_u De \alpha_{xx} + b_v De \alpha_{yy} - \left(1 + mqDe\sqrt{a^2 + b_u b_v} \right) \alpha_{xy} \right] \phi^q = 0$$

$$\left[2b_u De \alpha_{xy} - \left(1 + 2aDe + mqDe\sqrt{a^2 + b_u b_v} \right) \alpha_{yy} \right] \phi^q = 0$$

$$\alpha_{xx} = \frac{2b_v De \alpha_{xy}}{-1 + 2aDe - mqDe\sqrt{a^2 + b_u b_v}}$$

$$\alpha_{yy} = \frac{2b_u De \alpha_{xy}}{-1 + 2aDe + mqDe\sqrt{a^2 + b_u b_v}}$$

HOMOGENEOUS SOLUTION (3)

Back-substituting, three possible values of q and three possible stress fields

$$q = \frac{2}{m} - \frac{1}{mDe\sqrt{a^2 + b_u b_v}} \quad (1)$$



$$\alpha_{xx} = \frac{b_v \alpha_{xy}}{-a + \sqrt{a^2 + b_u b_v}}$$

$$\alpha_{yy} = \frac{b_u \alpha_{xy}}{a + \sqrt{a^2 + b_u b_v}}$$

$$q = -\frac{1}{mDe\sqrt{a^2 + b_u b_v}} \quad (2)$$



$$\alpha_{xx} = \frac{-b_v \alpha_{xy}}{a}$$

$$\alpha_{yy} = \frac{b_u \alpha_{xy}}{a}$$

$$q = -\frac{2}{m} - \frac{1}{mDe\sqrt{a^2 + b_u b_v}} \quad (3)$$



$$\alpha_{xx} = \frac{-b_v \alpha_{xy}}{a + \sqrt{a^2 + b_u b_v}}$$

$$\alpha_{yy} = \frac{b_u \alpha_{xy}}{a - \sqrt{a^2 + b_u b_v}}$$

Homogeneous solution is sum of all

Momentum not yet enforced

No boundary conditions imposed

MOMENTUM EQUATION (I)

$$\frac{\partial}{\partial y} \left(-\frac{\partial p}{\partial x} + \frac{\partial \tau_{xx}}{\partial x} + \frac{\partial \tau_{xy}}{\partial y} \right) - \frac{\partial}{\partial x} \left(-\frac{\partial p}{\partial y} + \frac{\partial \tau_{xy}}{\partial x} + \frac{\partial \tau_{yy}}{\partial y} \right) = 0$$

Case I

$m = -1$ avoids singularity at $x=0, y=0$ when $De \ll 1$

~~$b_u = \frac{a^2}{b_v}$~~ → singularities at all De

$$b_u = b_v$$

$$b_u = \frac{1 - 9a^2 De^2}{9b_v De^2}$$

$$b_u = \frac{1 - 4a^2 De^2}{4b_v De^2}$$

Stagnation + “vortex” flow

$$u = ax + b_u y$$

$$v = -ay + b_v x$$

possible forms to obey simultaneously momentum & UCM

Note change of signs at $De = 1/(3a)$ and $1/(2a)$

MOMENTUM EQUATION (2)

Case 2

$$\cancel{b_u = \frac{a^2}{b_v}} \longrightarrow \text{singularities}$$

$$b_u = \frac{1 - a^2 De}{b_v De^2} \longrightarrow \text{ok, but needs to be compatible with case (1)} \\ \text{(restrict values of } a \text{ and } De)$$

$$\cancel{b_u = -2ia - b_v}$$

$$\cancel{b_u = 2ia - b_v}$$

We will consider no contributions from case 2
to the solution ($k=0$ and $\alpha_{xy}=0$)

neglected at this stage

Case 3

After substitution of stresses all terms in equation are multiplied by $(1+m)$.
Since $m=-1$, momentum is automatically satisfied

STRESS FIELD

$$\tau_{xx} = \frac{b_v \alpha_{xy1}}{-a + \sqrt{a^2 + b_u b_v}} \phi^{\frac{2}{m} - \frac{1}{mDe\sqrt{a^2 + b_u b_v}}} - \frac{b_v \alpha_{xy3}}{a + \sqrt{a^2 + b_u b_v}} \phi^{-\frac{2}{m} - \frac{1}{mDe\sqrt{a^2 + b_u b_v}}} - \frac{2(a + 2Dea^2 + b_u b_v De + b_v^2 De)}{4a^2 De^2 - 1 + 4b_u b_v De^2}$$

$$\tau_{xy} = \alpha_{xy1} \phi^{\frac{2}{m} - \frac{1}{mDe\sqrt{a^2 + b_u b_v}}} + \alpha_{xy3} \phi^{-\frac{2}{m} - \frac{1}{mDe\sqrt{a^2 + b_u b_v}}} - \frac{b_u + b_v + 2aDe(b_u - b_v)}{4a^2 De^2 - 1 + 4b_u b_v De^2}$$

$$\tau_{yy} = \frac{b_u \alpha_{xy1}}{a + \sqrt{a^2 + b_u b_v}} \phi^{\frac{2}{m} - \frac{1}{mDe\sqrt{a^2 + b_u b_v}}} + \frac{b_u \alpha_{xy3}}{a - \sqrt{a^2 + b_u b_v}} \phi^{-\frac{2}{m} - \frac{1}{mDe\sqrt{a^2 + b_u b_v}}} - \frac{2(-a + 2Dea^2 + b_u b_v De + b_u^2 De)}{4a^2 De^2 - 1 + 4b_u b_v De^2}$$

with $\alpha_{xy1} = \alpha_{xy1}(a, b_u, b_v)$, $\alpha_{xy3} = \alpha_{xy3}(a, b_u, b_v)$ such as

$$\alpha_{xy1} = \frac{\alpha_1(-a + \sqrt{a^2 + b_u b_v})}{b_v} \quad \alpha_{xy3} = \frac{\alpha_3(a + \sqrt{a^2 + b_u b_v})}{b_v}$$

$a = 1, b_u = 0, b_v = 0 \Rightarrow$ Becherer et al. JNNFM 153 (2008) 183

STREAMLINES AND STRESSES (I)

Stream function $\psi = axy + b_u \frac{y^2}{2} - b_v \frac{x^2}{2}$

Stream function of vortex $\psi_1 = \psi_{total} - \psi_{stagnation} = b_u \frac{y^2}{2} - b_v \frac{x^2}{2}$

$$b_u = b_v$$

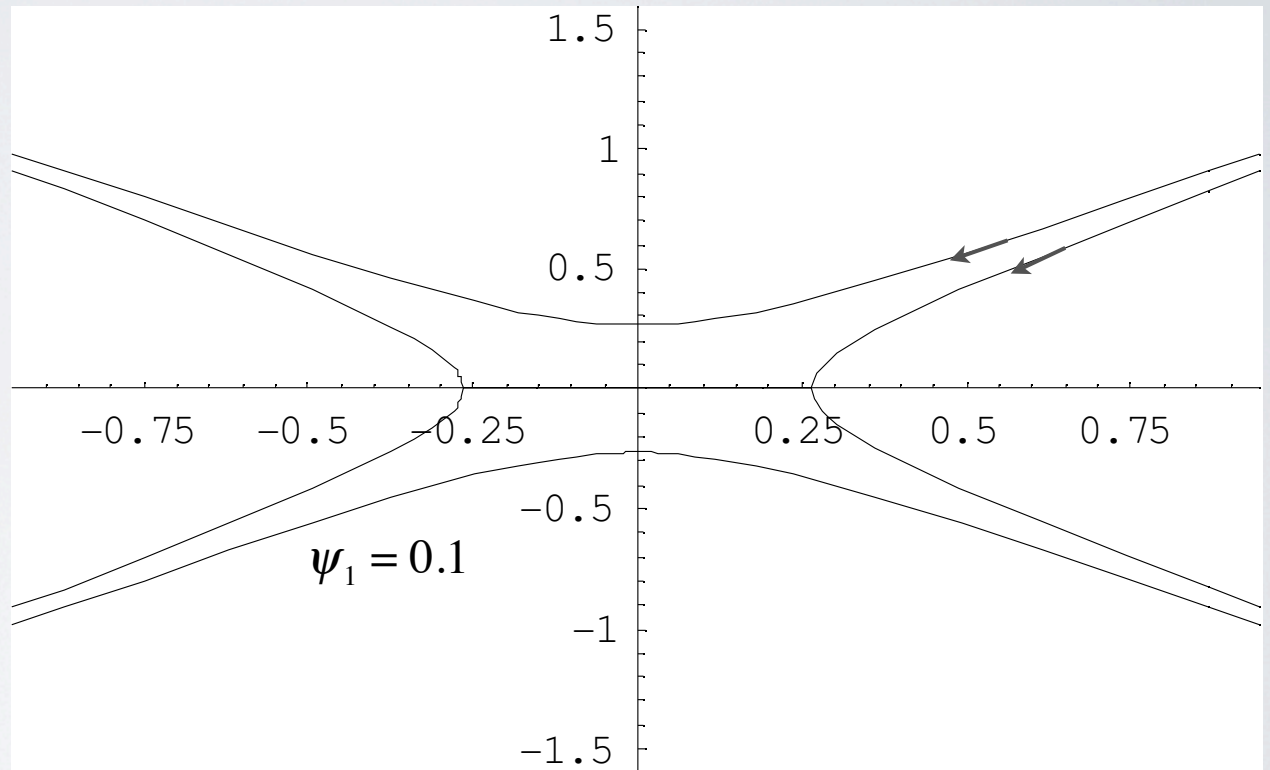
$$b_u = \frac{1 - 9a^2 De}{9b_v De^2}$$

$$b_u = \frac{1 - 4a^2 De}{4b_v De^2}$$

(I) $De < \frac{1}{\sqrt{9a^2}}$

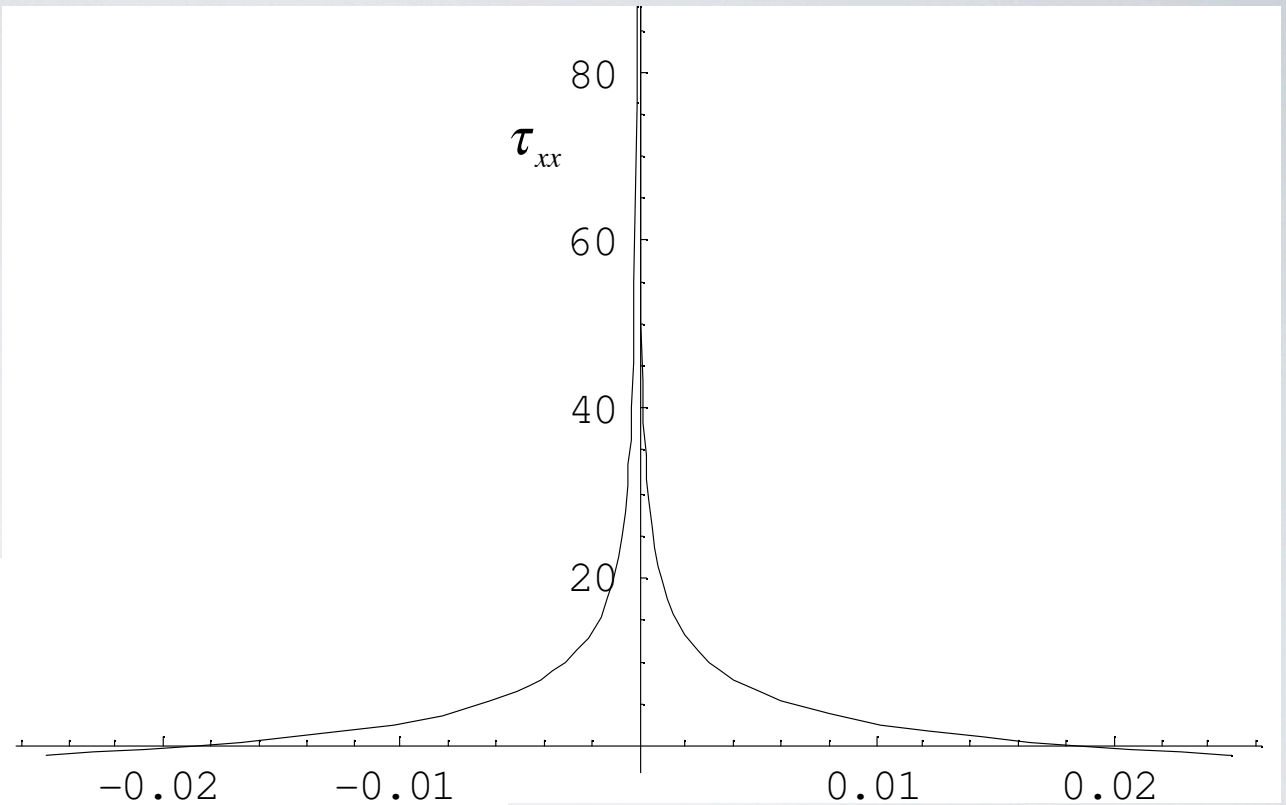
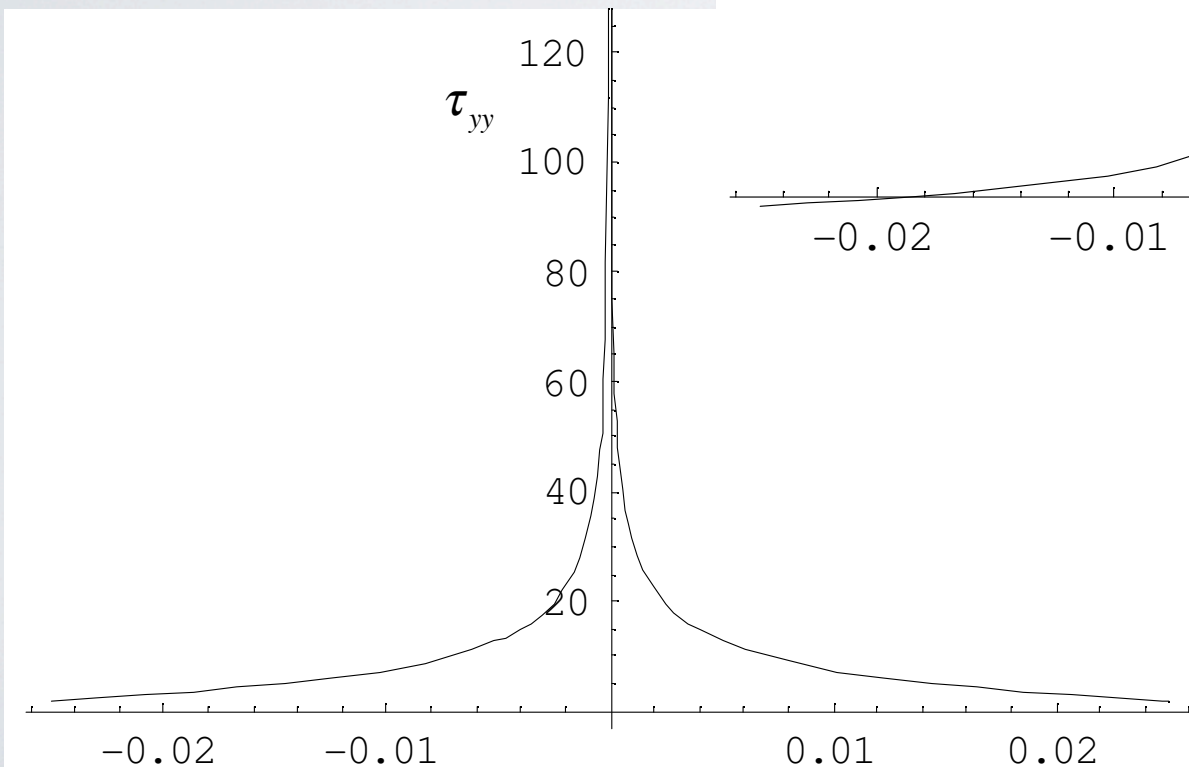
$$b_u = b_v = 2.85; a = -1;$$

$$De = 0.2$$



STREAMLINES AND STRESSES (2)

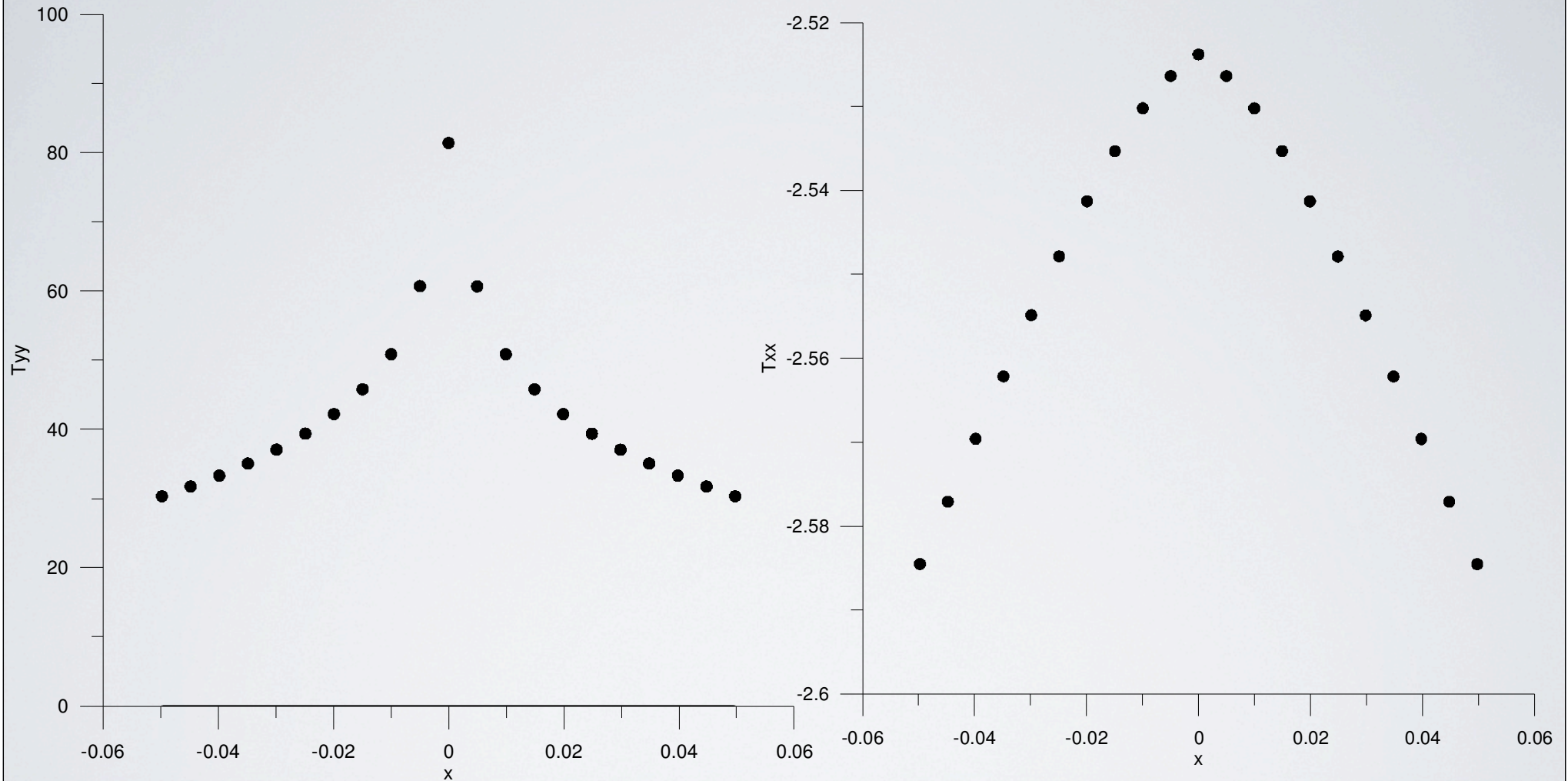
$$y = 0.000025$$



Transitions in some stagnation viscoelastic flows at $Re=0$
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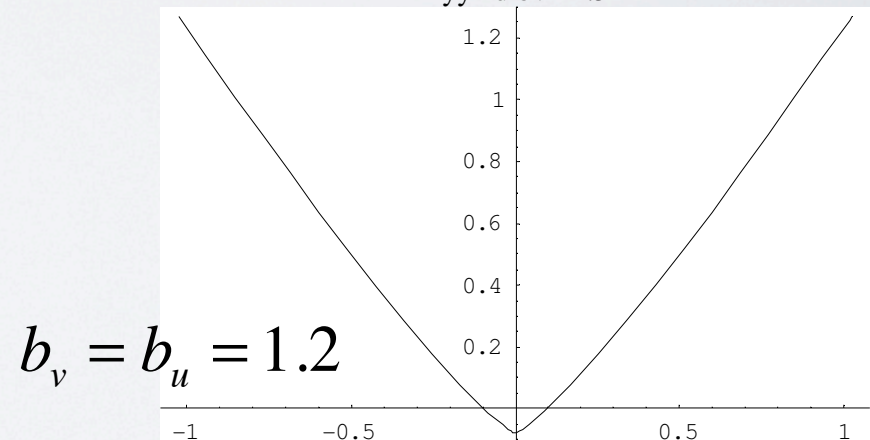
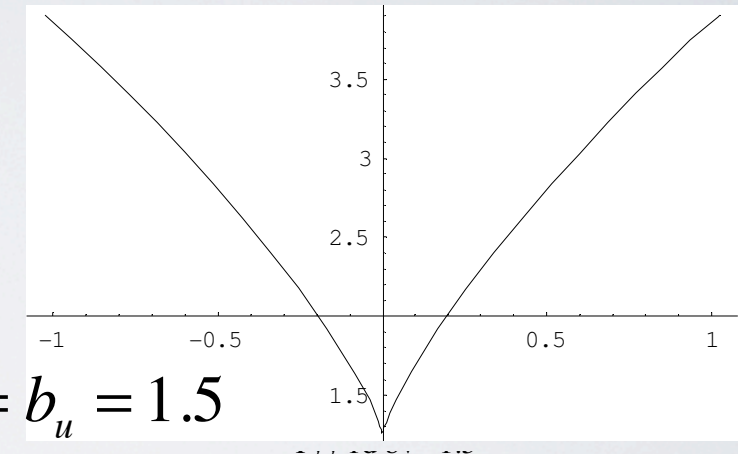
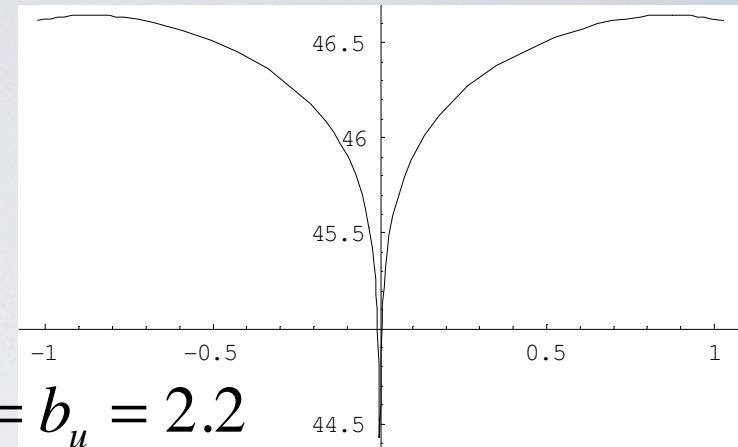
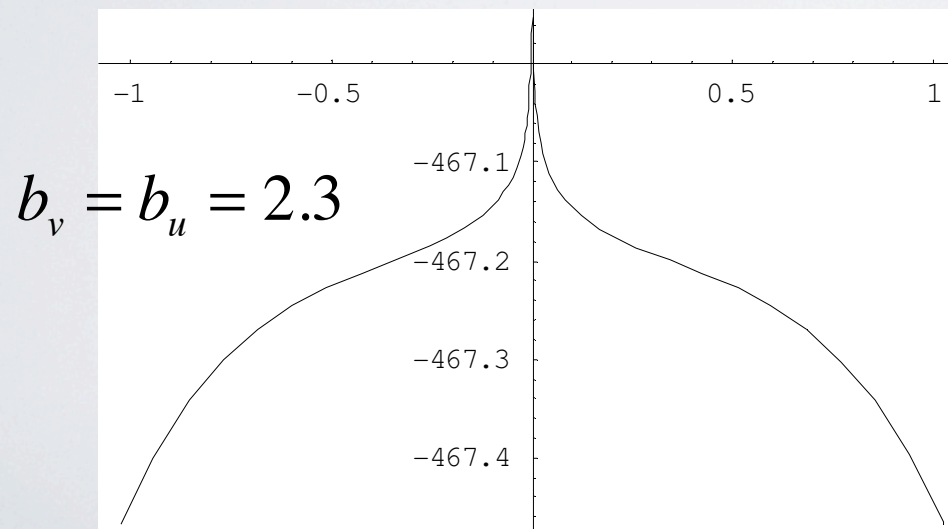
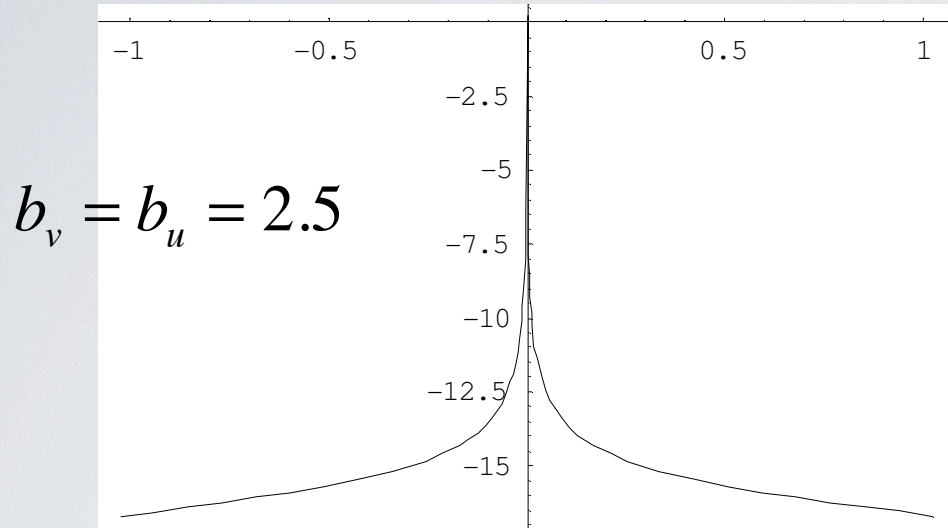
STRESSES FROM NUMERICAL RESULTS: DE=0.2 (I)



STRESSES (2)

(Ia) $De < \frac{1}{\sqrt{9a^2}}$ $De = 0.2; a = -1$

τ_{xx}



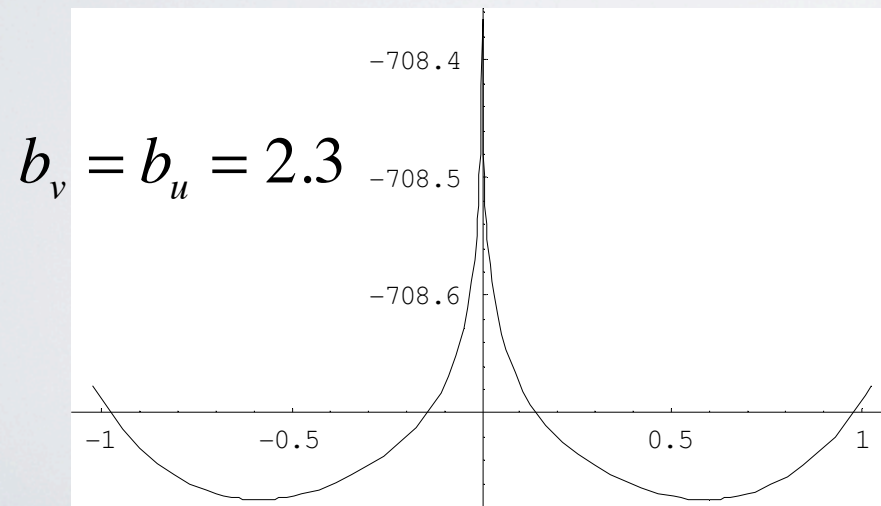
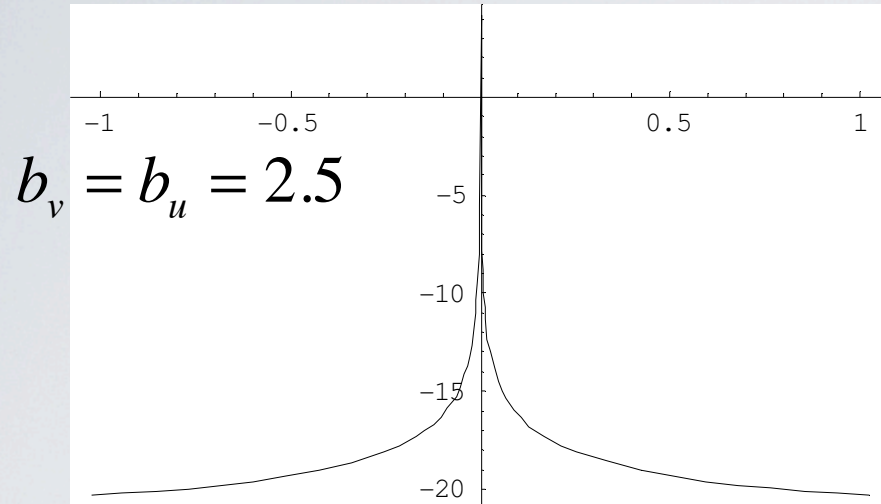
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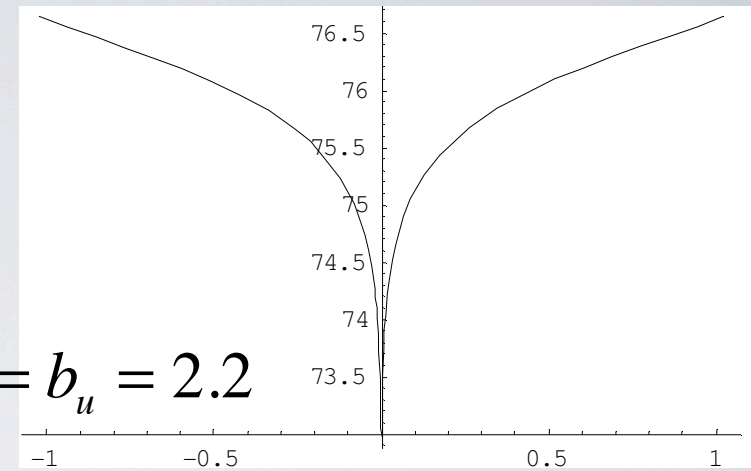
STRESSES (3)

(I) $De < \frac{1}{\sqrt{9a^2}}$ $De = 0.2; a = -1$

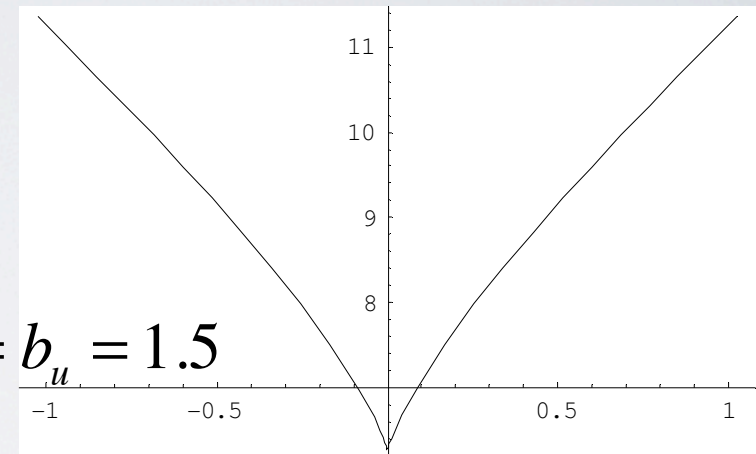
τ_{yy}



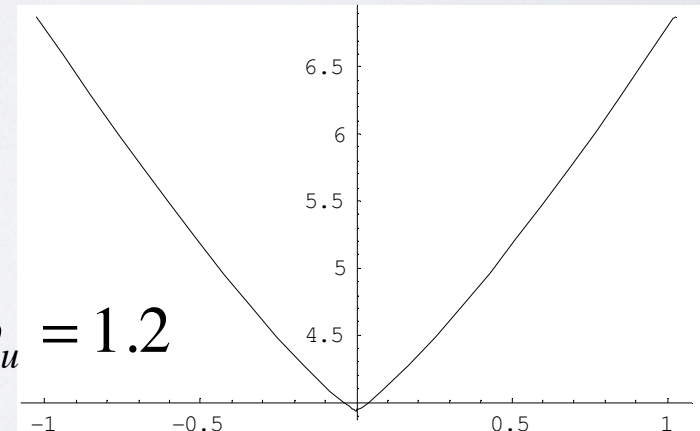
$b_v = b_u = 2.2$



$b_v = b_u = 1.5$



$b_v = b_u = 1.2$



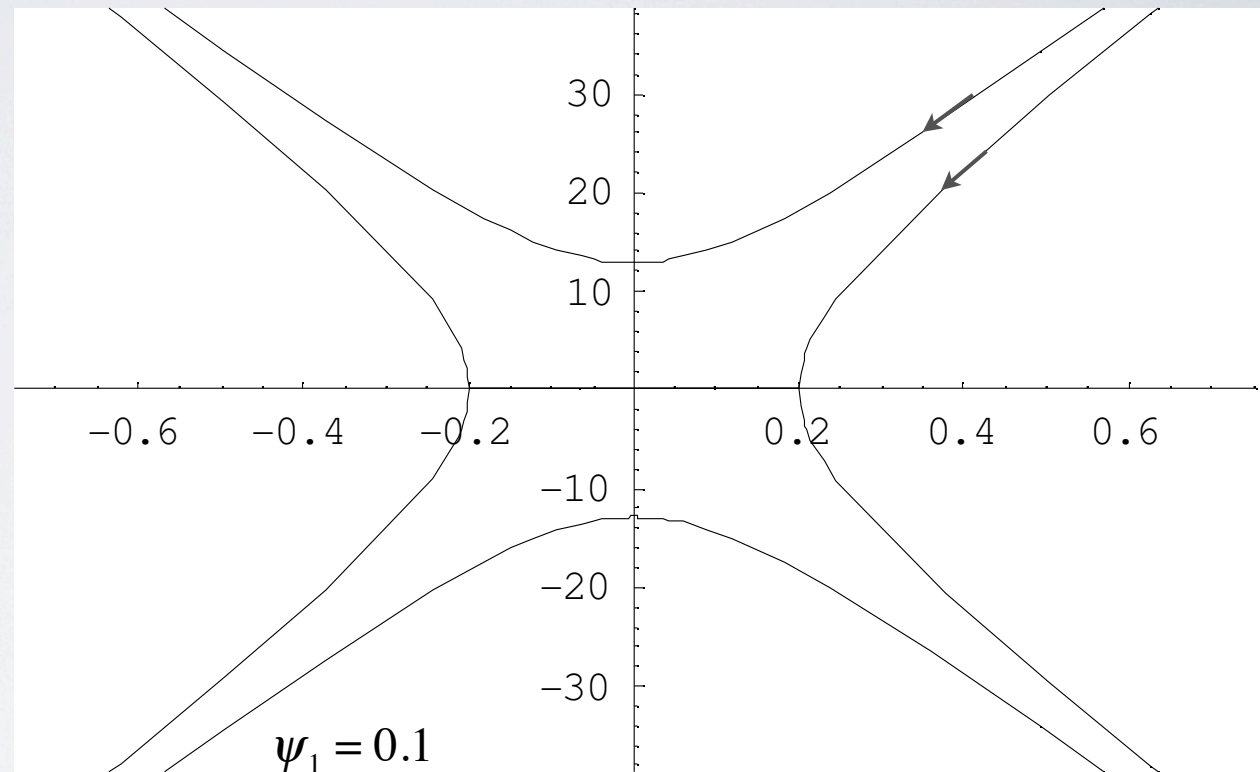
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STREAMLINES AND STRESSES (3)

$$(2) \quad De < \frac{1}{\sqrt{9a^2}}$$

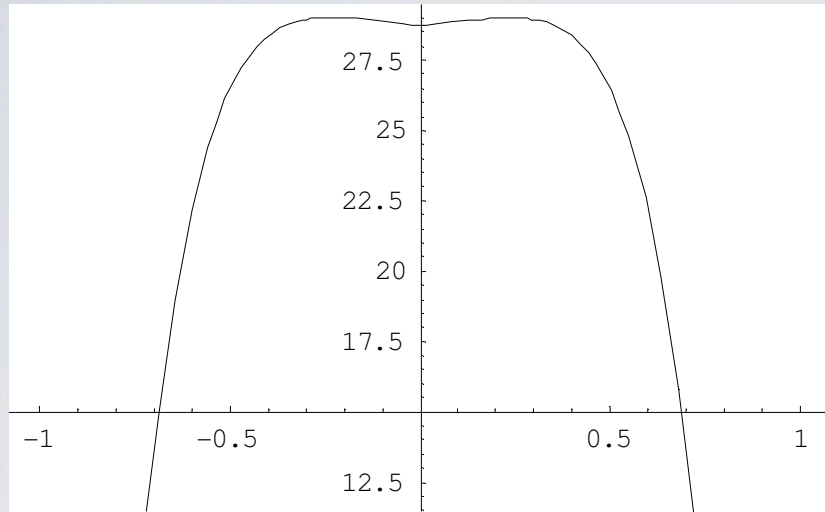
$$b_u = \frac{1 - 9a^2 De^2}{9b_v De^2}; a = -1; b_v = 5; De = \frac{1}{3\sqrt{a^2}} - 0.001$$



STREAMLINES AND STRESSES (4)

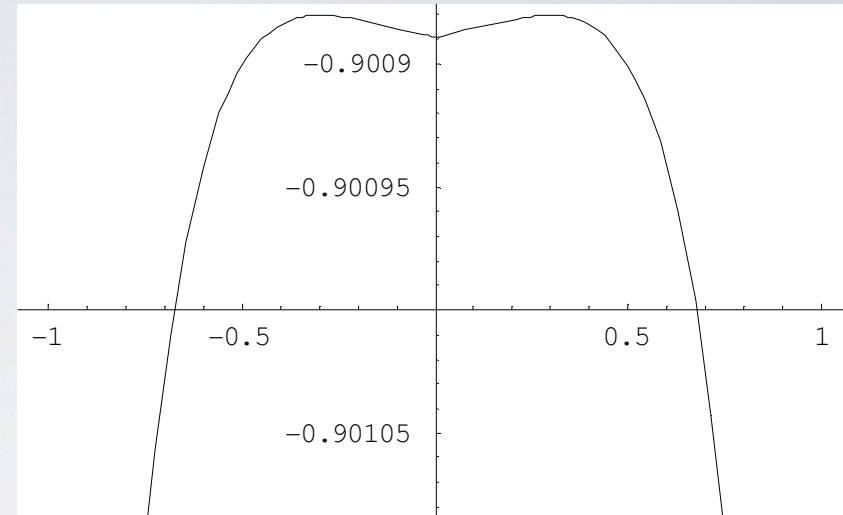
$b_v = 5$

τ_{xx}

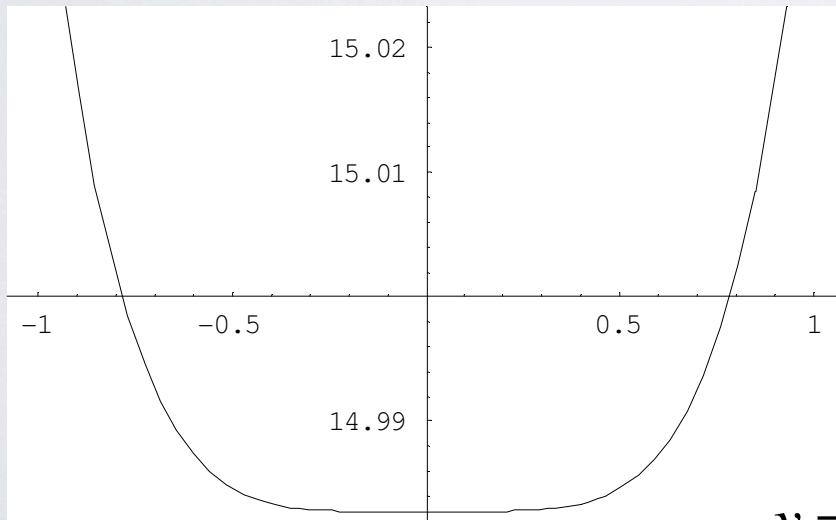


τ_{xx}

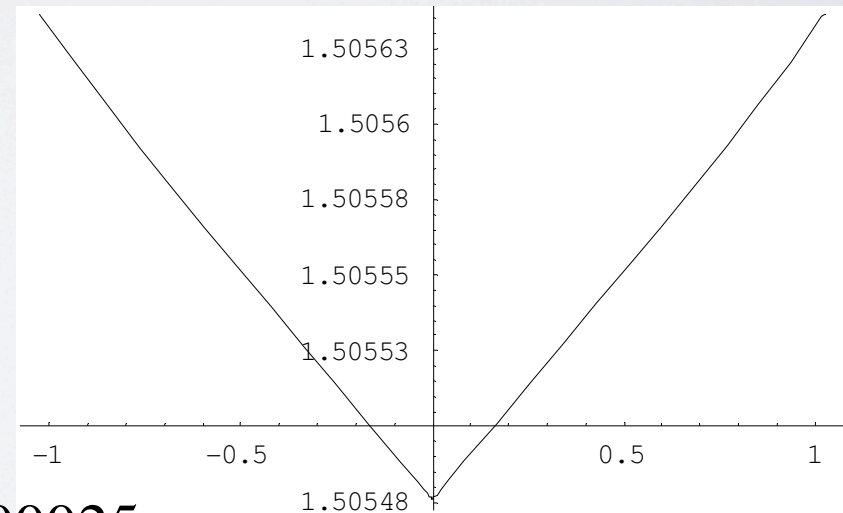
$b_v = 0.5$



τ_{yy}



τ_{yy}



$y = 0.000025$

STREAMLINES AND STRESSES (5)

$$(3) De > \frac{1}{\sqrt{9a^2}}$$

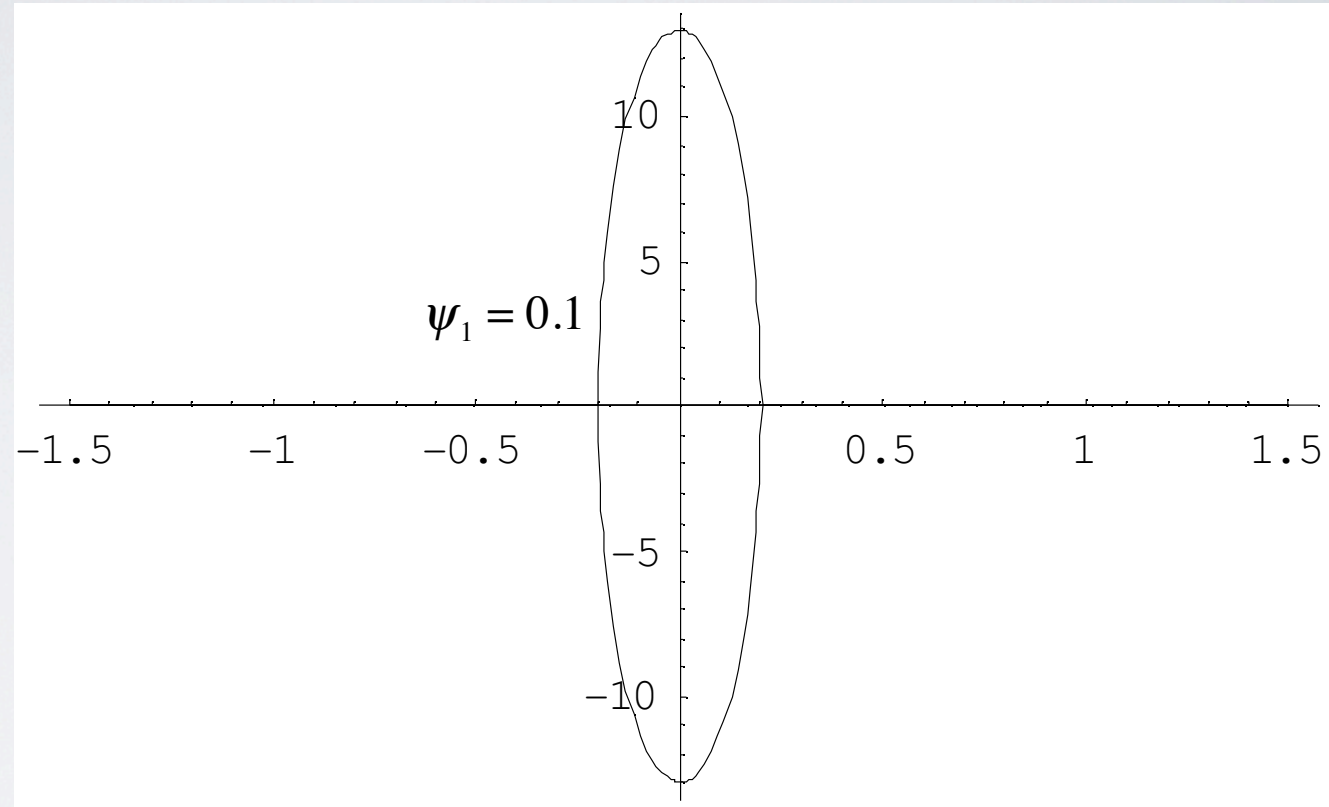
$$b_u = \frac{1 - 9a^2 De^2}{9b_v De^2}; a = -1; b_v = 5; De = \frac{1}{3\sqrt{a^2}} + 0.001$$

b_u, b_v

Opposite signs



Vortex



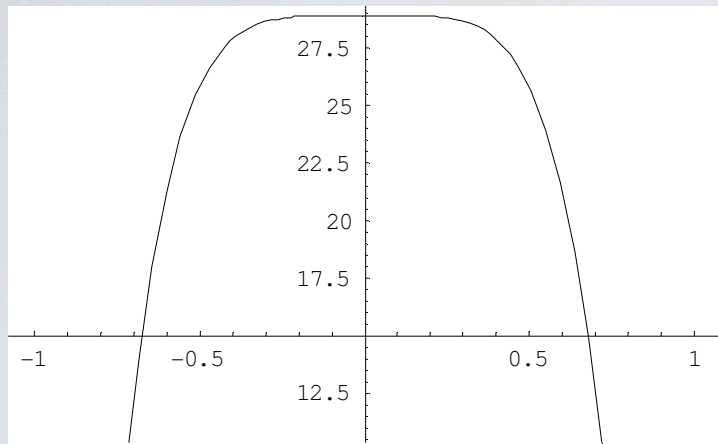
$De < De_c$ **Vortex enclosing stagnation point is not possible.**

$De > De_c$ **Vortex enclosing stagnation point is possible.**

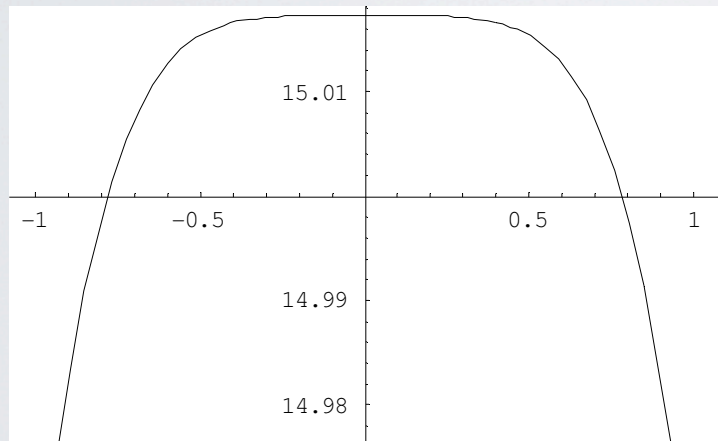
STREAMLINES AND STRESSES (6)

$$b_v = 5$$

$$\tau_{xx}$$

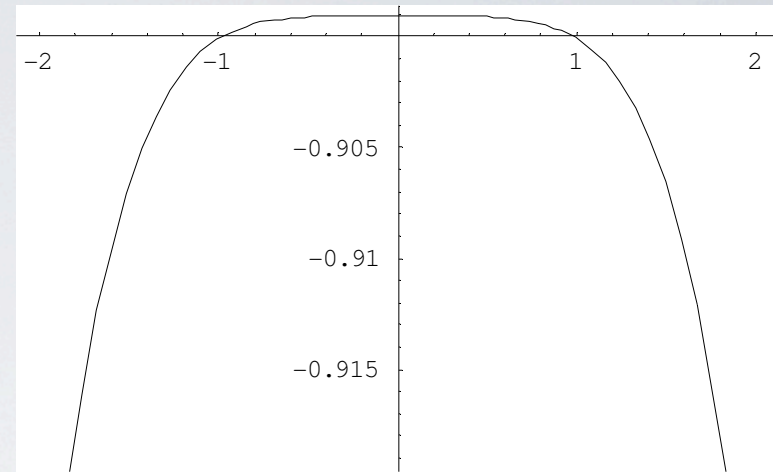


$$\tau_{yy}$$

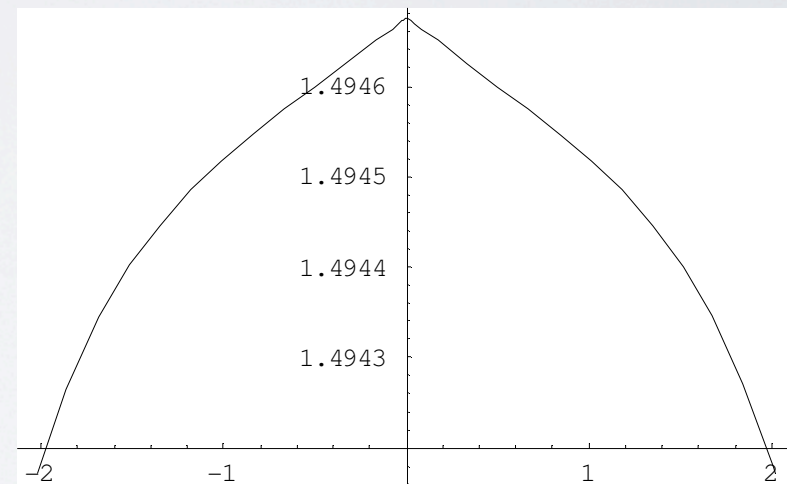


$$\tau_{xx}$$

$$b_v = 0.5$$



$$\tau_{yy}$$

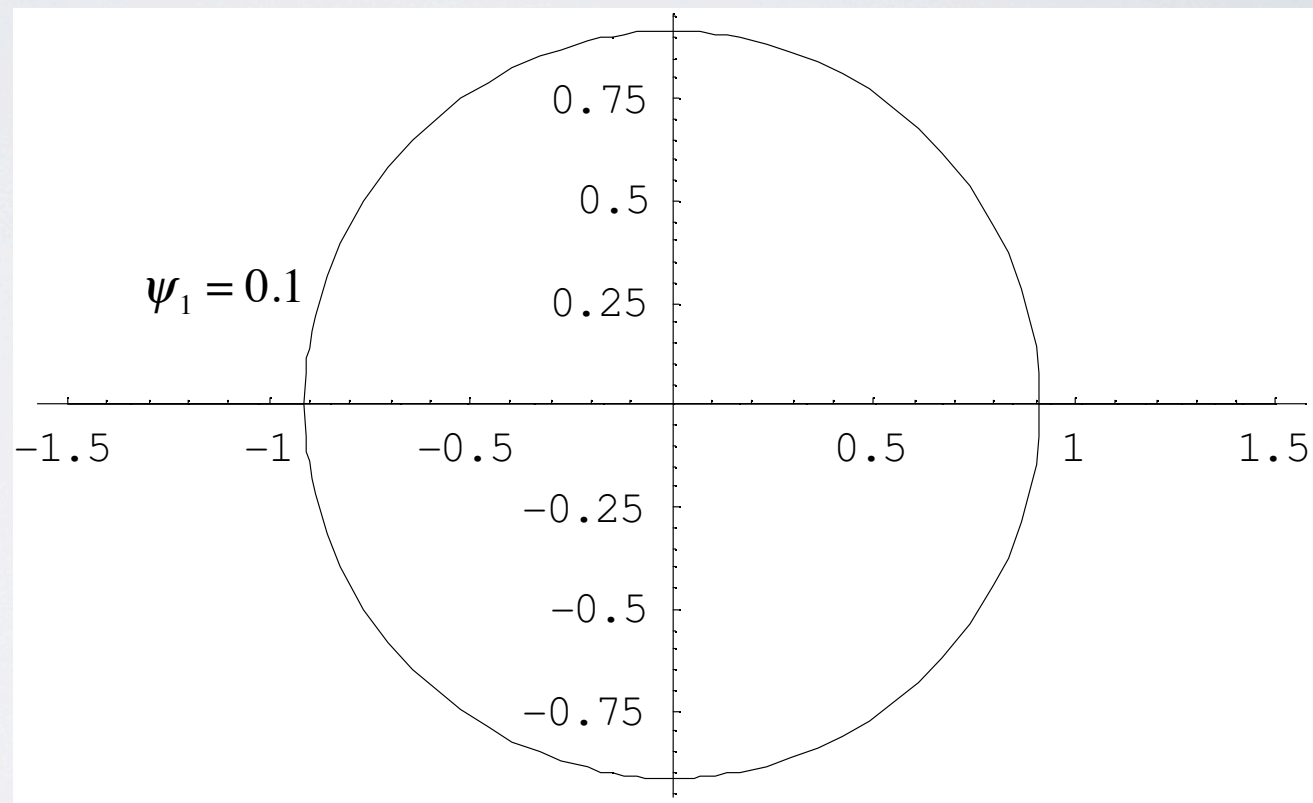


$$y = 0.000025$$

STREAMLINES AND STRESSES (7)

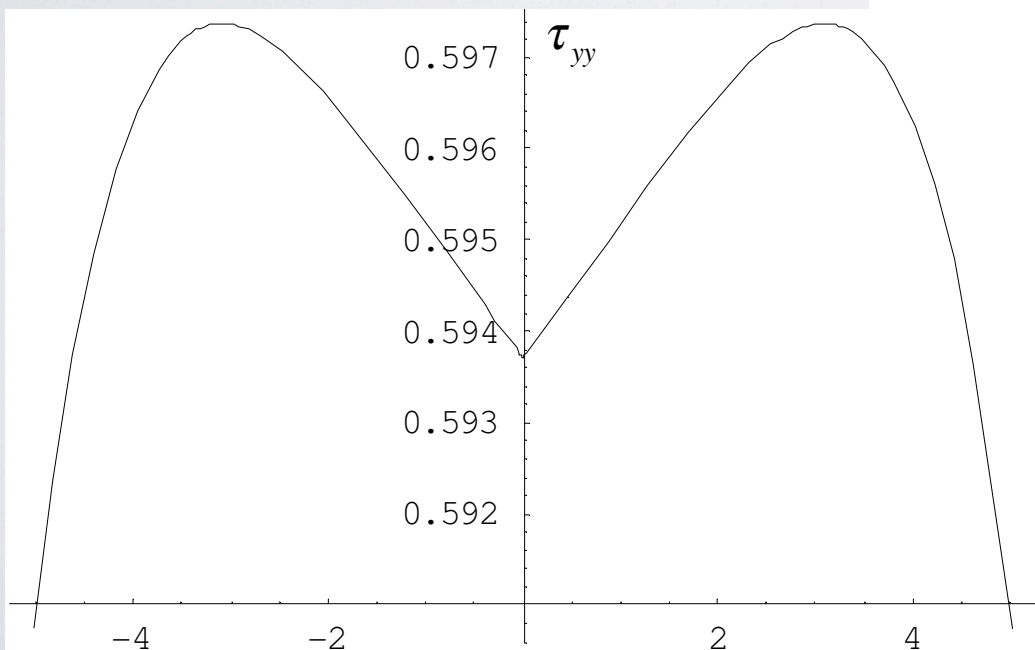
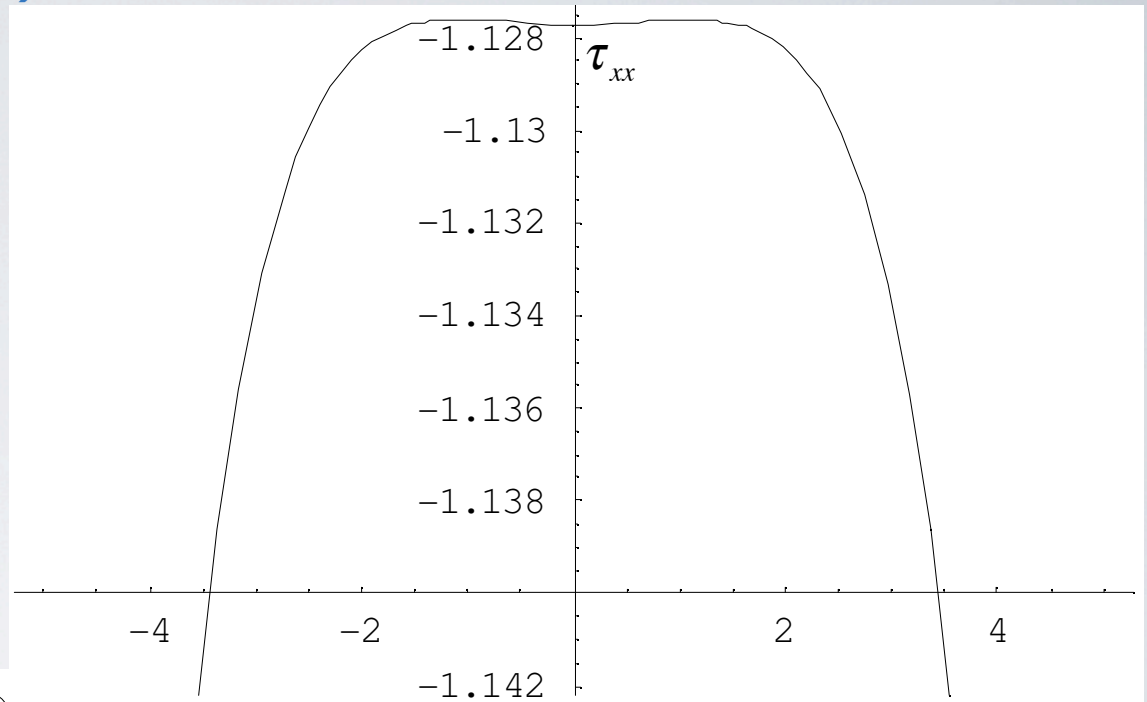
(4) $De > \frac{1}{\sqrt{9a^2}}$ - Circular vortex

$$b_u = \frac{1 - 9a^2 De^2}{9b_v De^2}; a = -1; b_v = 0.23999; De = \frac{1}{3\sqrt{a^2}} + 0.01$$



STREAMLINES AND STRESSES (8)

$$y = 0.000025$$



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CLOSURE

- Steady symmetric to steady asymmetric is a purely elastic instability. Inertia and solvent delays and eliminates this transition.
- This transition exists with bounded extensional viscosity, but is weakened with ϵ
- Steady asymmetric flow is a combination of a planar stagnation and a vortex
- Analytical solution obtained enforcing UCM constitutive equation and momentum. It shows closed vortex cannot exist below $De < 1/(3a)$
- Behavior of the solution currently under investigation: need to impose BC
- Need for stability analysis on the analytical solution.

ACKNOWLEDGEMENTS

- Fundação para a Ciência e a Tecnologia & Feder: PTDC/EQU-FTT/71800/2006, PTDC/EQU-FTT/70727/2006, PTDC/EME-MFE/70186/2006
- CNpQ 200120/2009-3